

LANL, 24 April 2002

The Fate of the First Stars

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Overview

I Introduction

- Basics of stellar evolution
- The life of modern stars

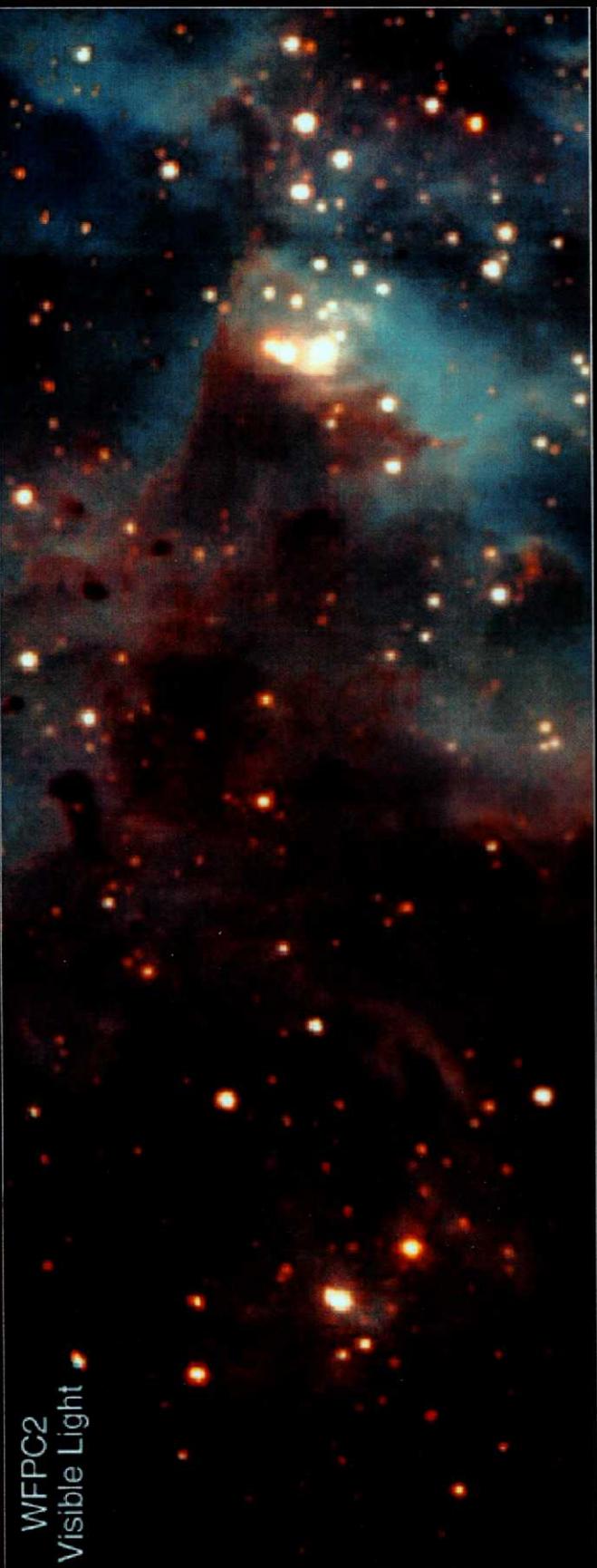
II Primordial Stars

- Birth
- Life and death
- The ashes

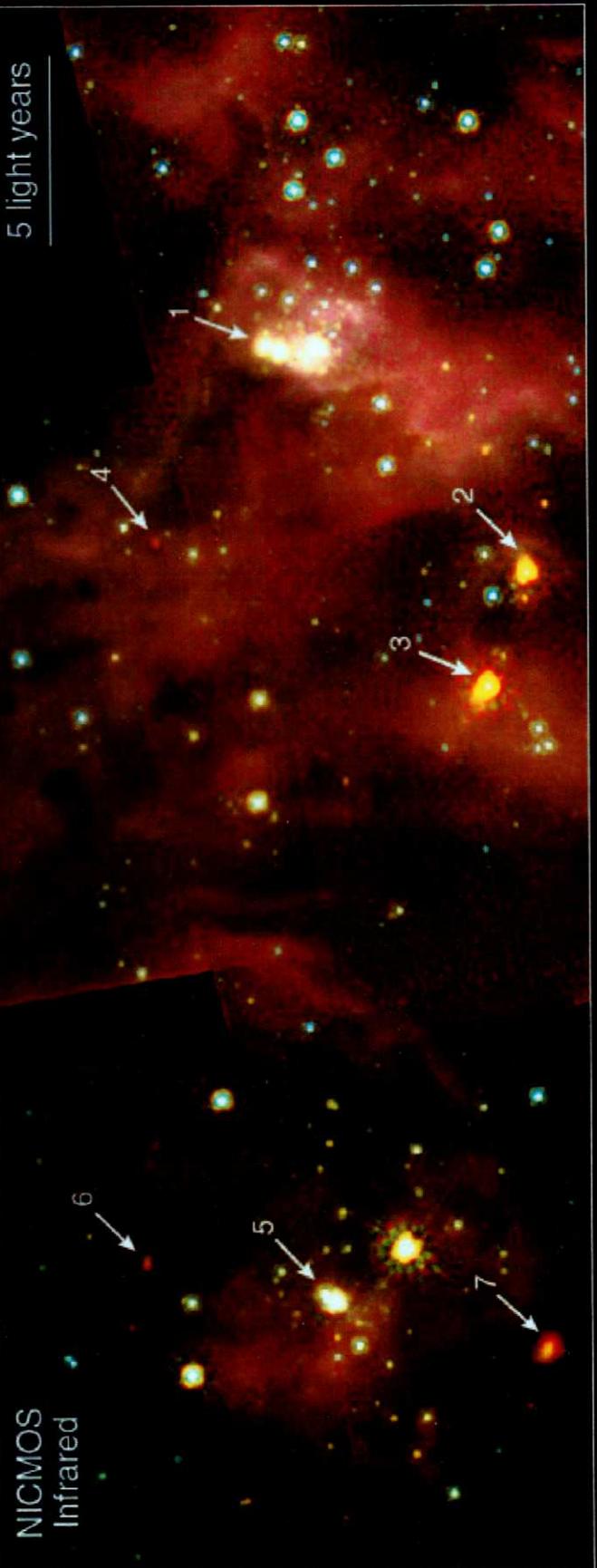
III Conclusions

<http://firststars.org>

WFPC2
Visible Light



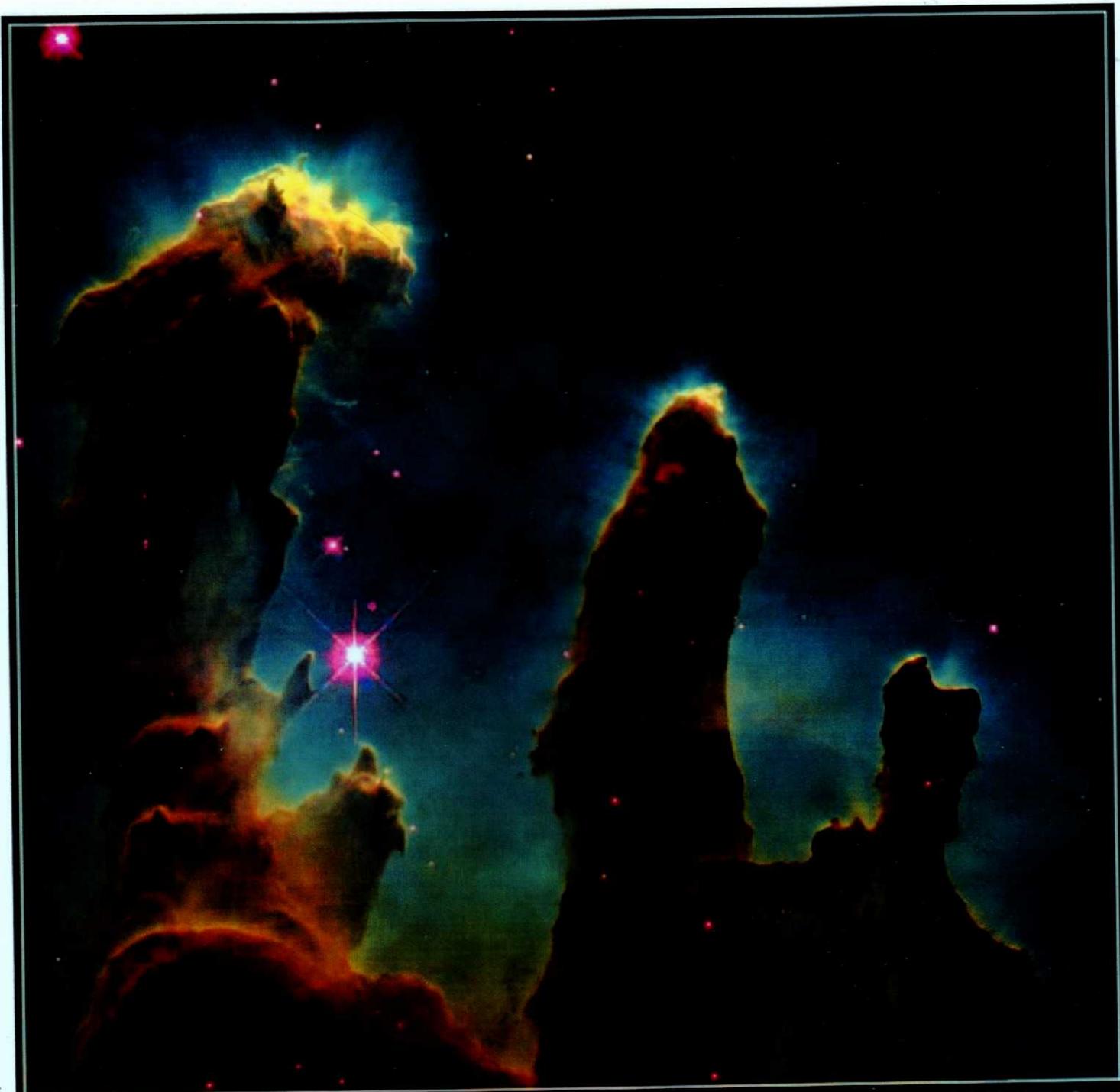
NICMOS
Infrared



30 Doradus Nebula Details

HST • WFPC2 • NICMOS

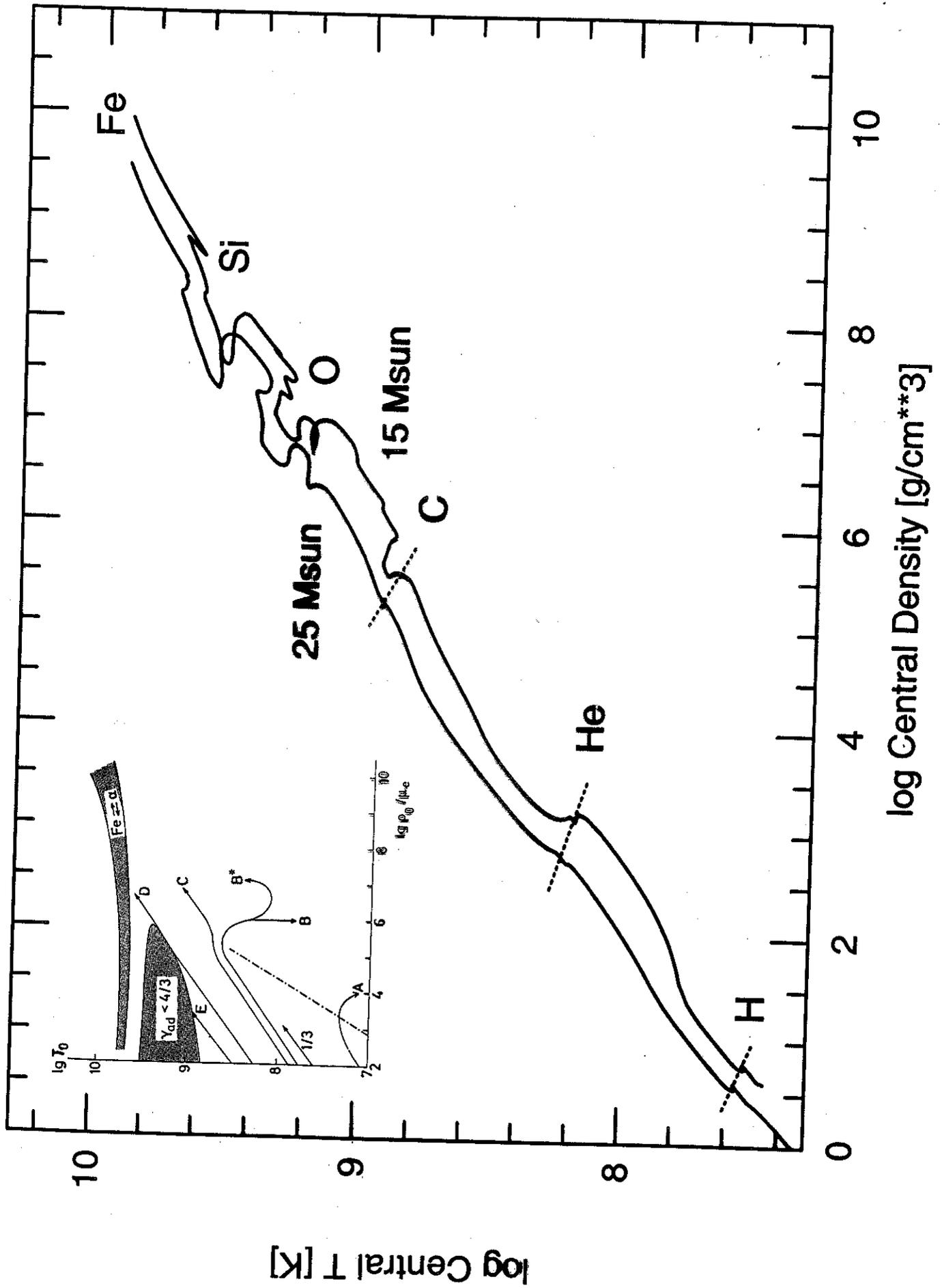
PRC99-33b • STScI OPO • N. Walborn (STScI), R. Barbá (La Plata Observatory) and NASA



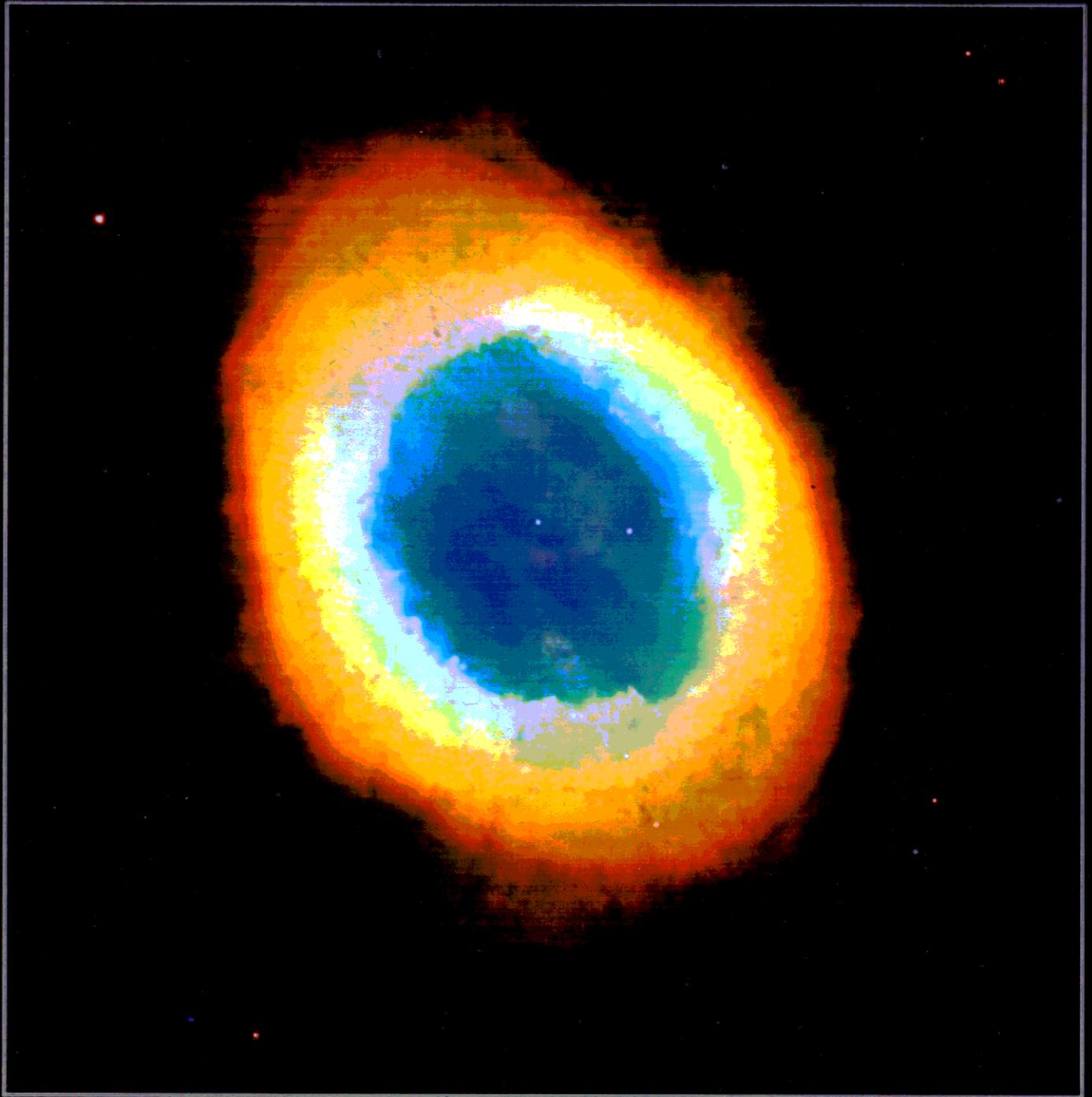
Gaseous Pillars • M16

HST • WFPC2

PRC95-44a • ST ScI OPO • November 2, 1995
J. Hester and P. Scowen (AZ State Univ.), NASA

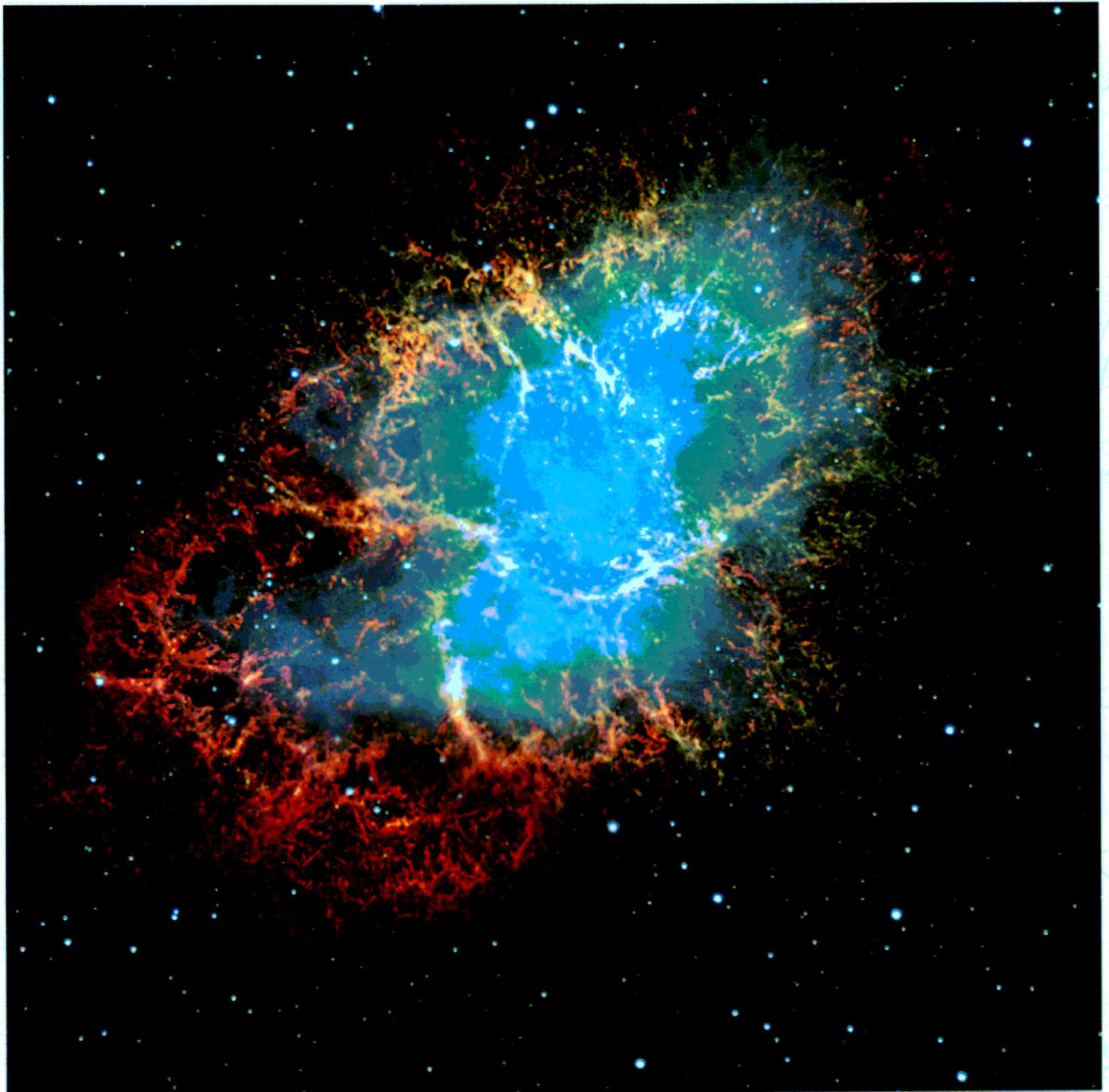


Ring Nebula



Hubble
Heritage

PRC99-01 • Space Telescope Science Institute • Hubble Heritage Team (AURA/STScI/NASA)



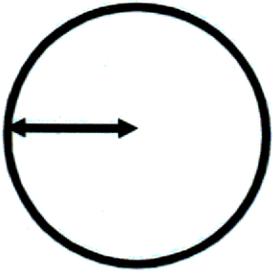
The Crab Nebula in Taurus (VLT KUEYEN + FORS2)

ESO PR Photo 40f/99 (17 November 1999)

© European Southern Observatory



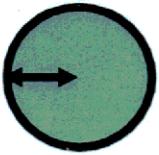
10,000 km



white dwarf

density: $1,000,000 \text{ g/cm}^3$

10 km



neutron star

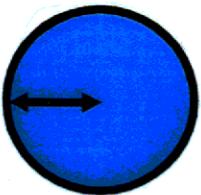
density: $\sim 100,000,000,000,000 \text{ g/cm}^3$

3 km



black hole

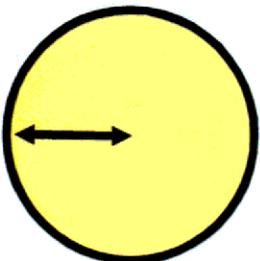
6,000 km



earth

density: 5.5 g/cm^3

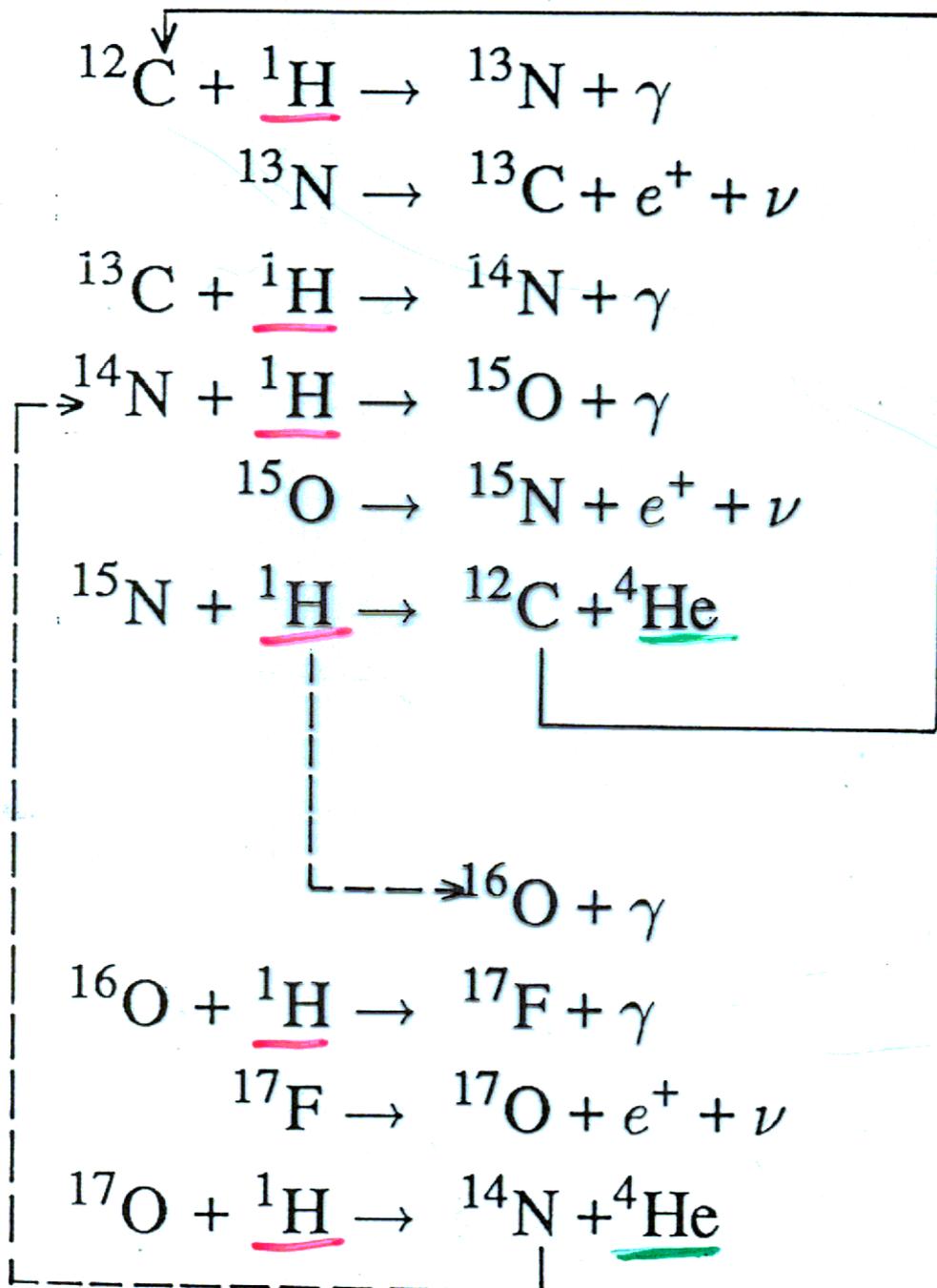
700,000 km



sun

density: 1.4 g/cm^3

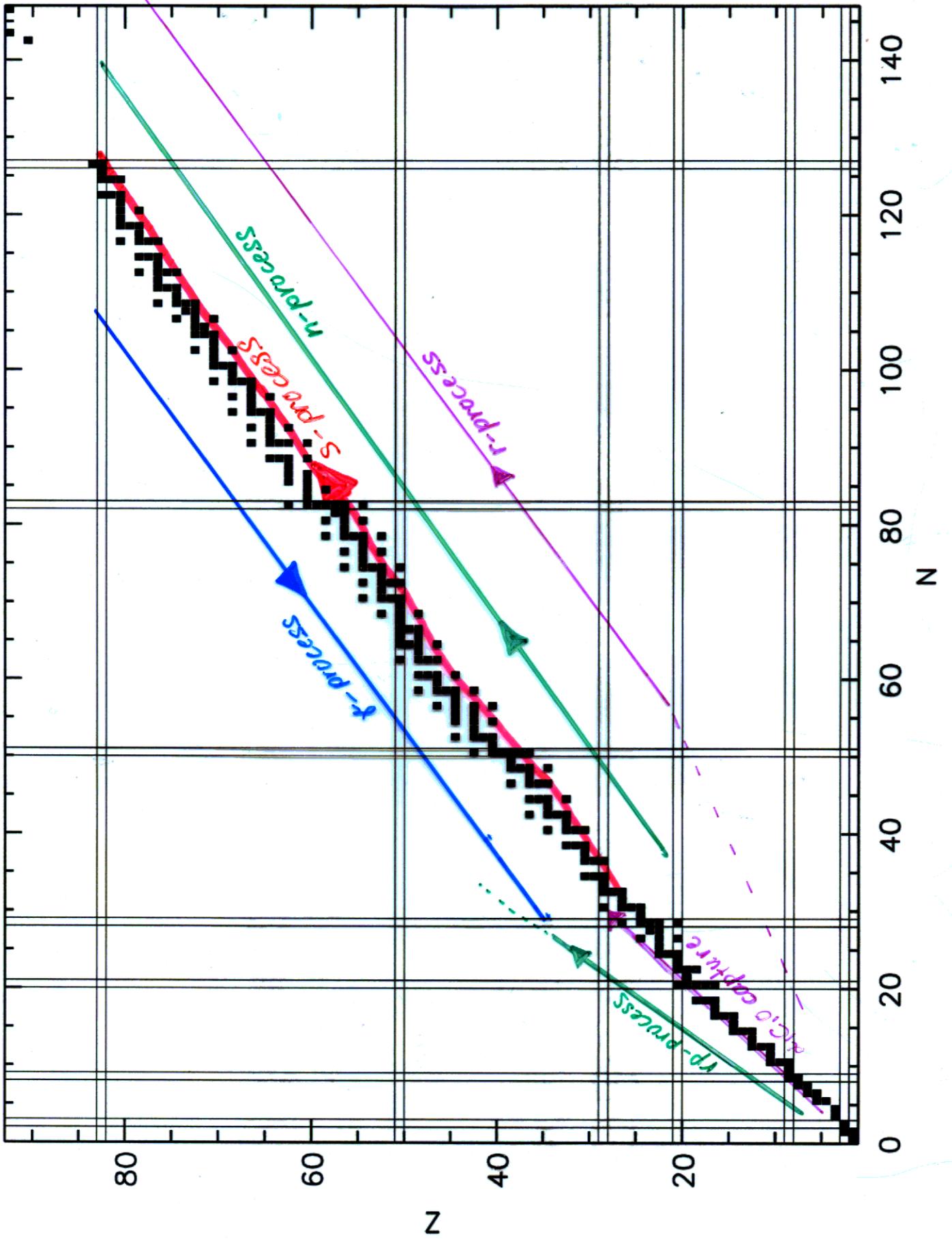
THE CNO_{Bi}-CYCLE



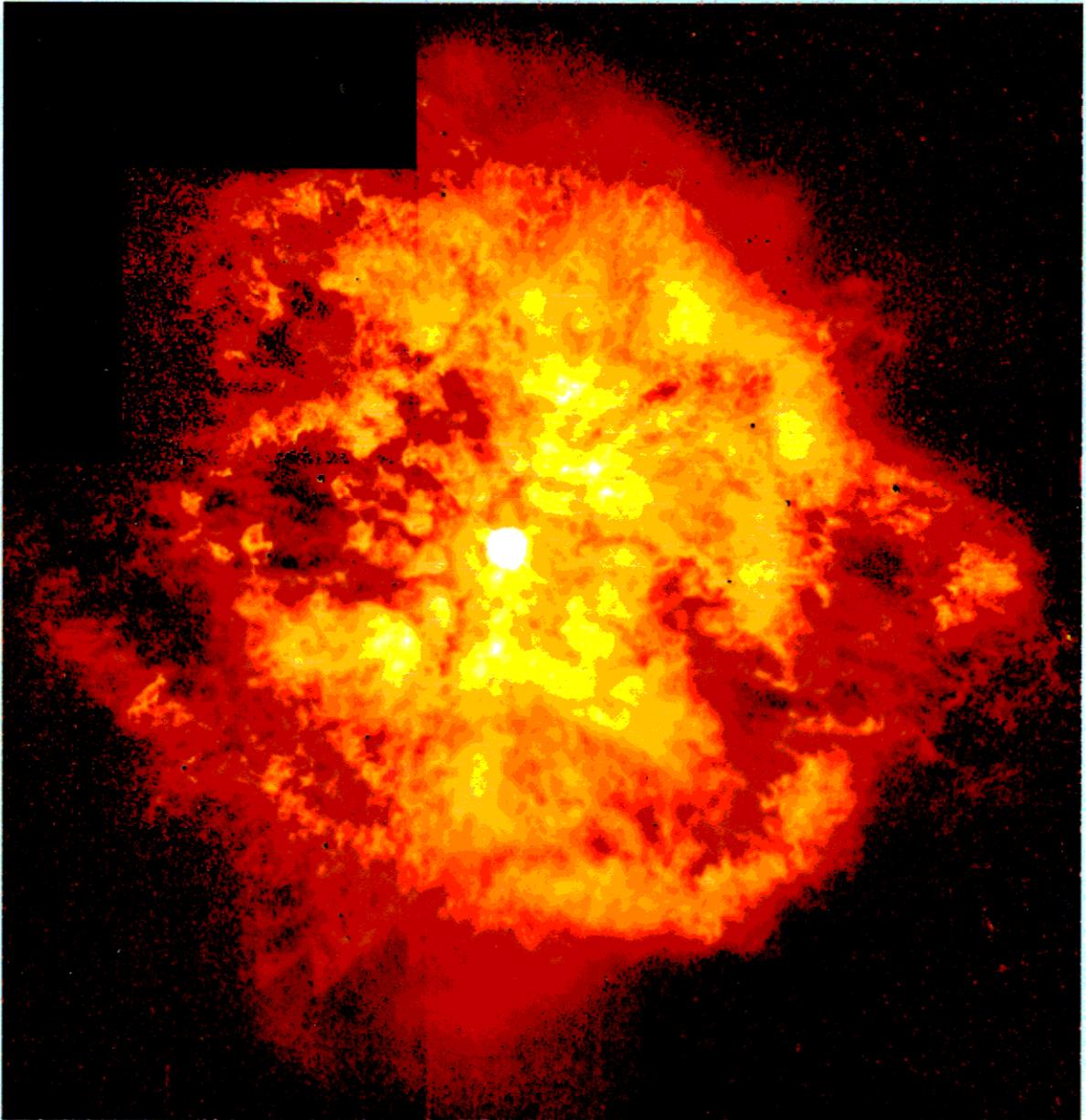
Advanced Nuclear Burning Stages

(e.g., 20 solar masses)

Fuel	Main Product	Secondary Products	Temp (10 ⁹ K)	Time (yr)	Time
H	He	¹⁴ N	0.02	10 ⁷	$4H \rightarrow {}^4\text{He}$
He	C, O	¹⁸ O, ²² Ne s-process	0.2	10 ⁶	$3\alpha \rightarrow {}^{12}\text{C}$ ${}^{12}\text{C}(\alpha, \gamma){}^{16}\text{O}$
C	Ne, Mg	Na	0.8	10 ³	${}^{12}\text{C} + {}^{12}\text{C}$
Ne	O, Mg	Al, P	1.5	3	${}^{20}\text{Ne}(\gamma, \alpha)$
O	Si, S	Cl, Ar K, Ca	2.0	0.8	${}^{16}\text{O} + {}^{16}\text{O}$
Si	Fe	Ti, V, Cr Mn, Co, Ni	3.5	1 week	${}^{28}\text{Si}(\gamma, \alpha)$



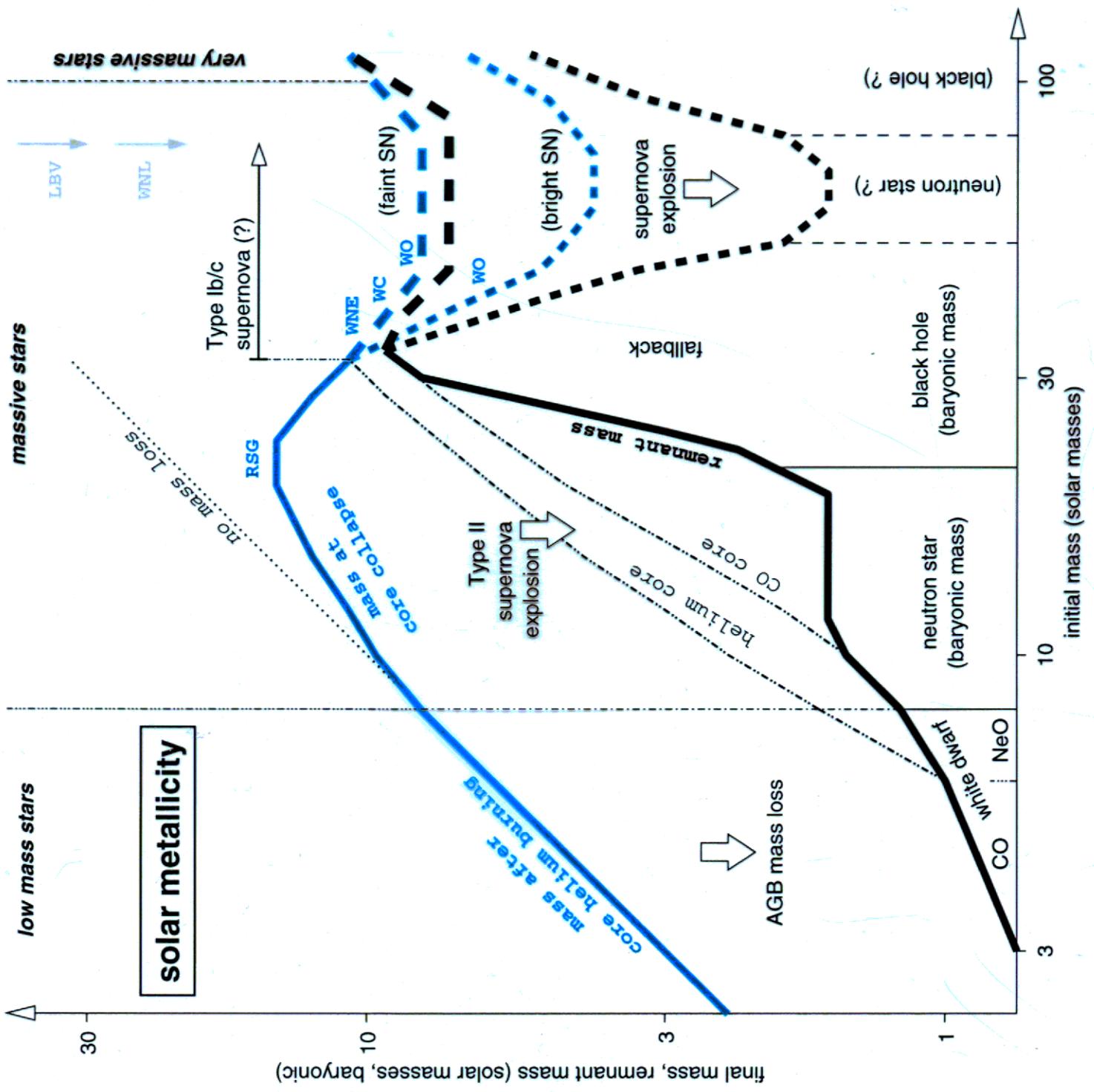
WR 124



STScI-PR98-38

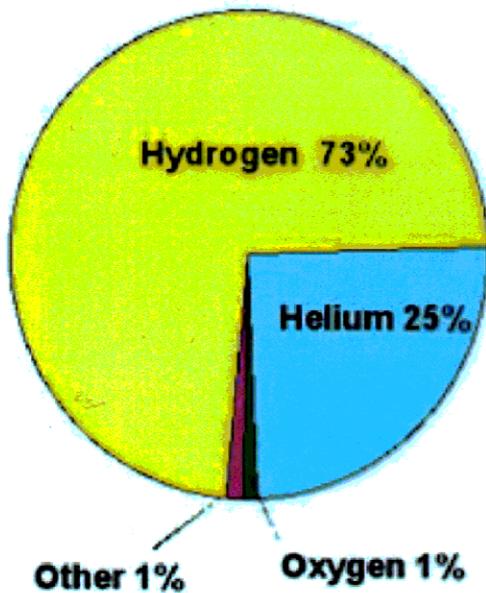
[http://OPOBITE.STSCI.EDU/
PUPINFO/PR/1998/38/PR-PHOTOS.HTML](http://OPOBITE.STSCI.EDU/PUPINFO/PR/1998/38/PR-PHOTOS.HTML)

WR 124 in M1-67
DIAMETER OF THE NEBULA \sim 1000 A.U.
VERY CLUMPY WIND

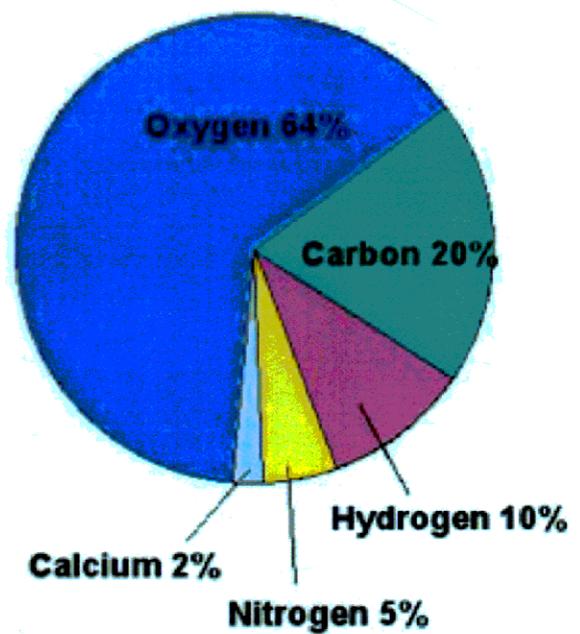


Relative Abundance by Weight

Universe



Humans



FORMATION OF A 200 M_{\odot} PROTOSTELLAR CORE @ $z=19.1$

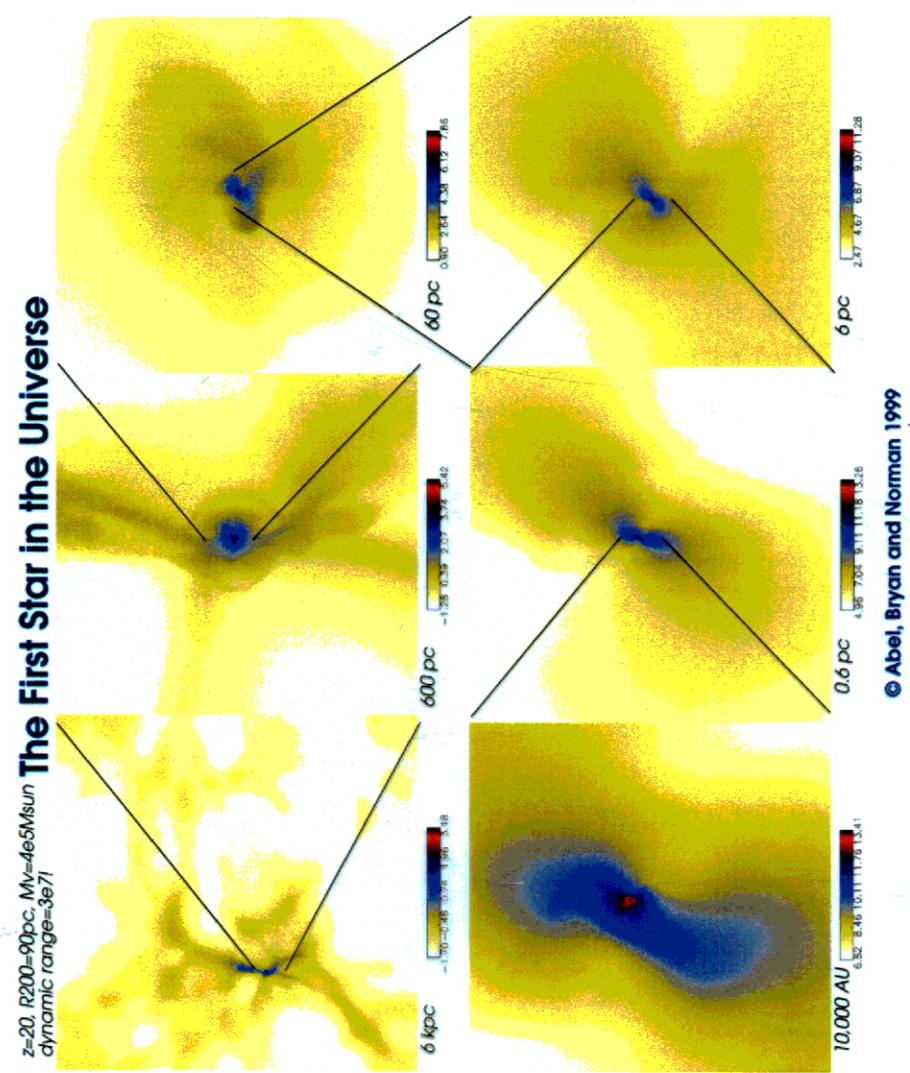


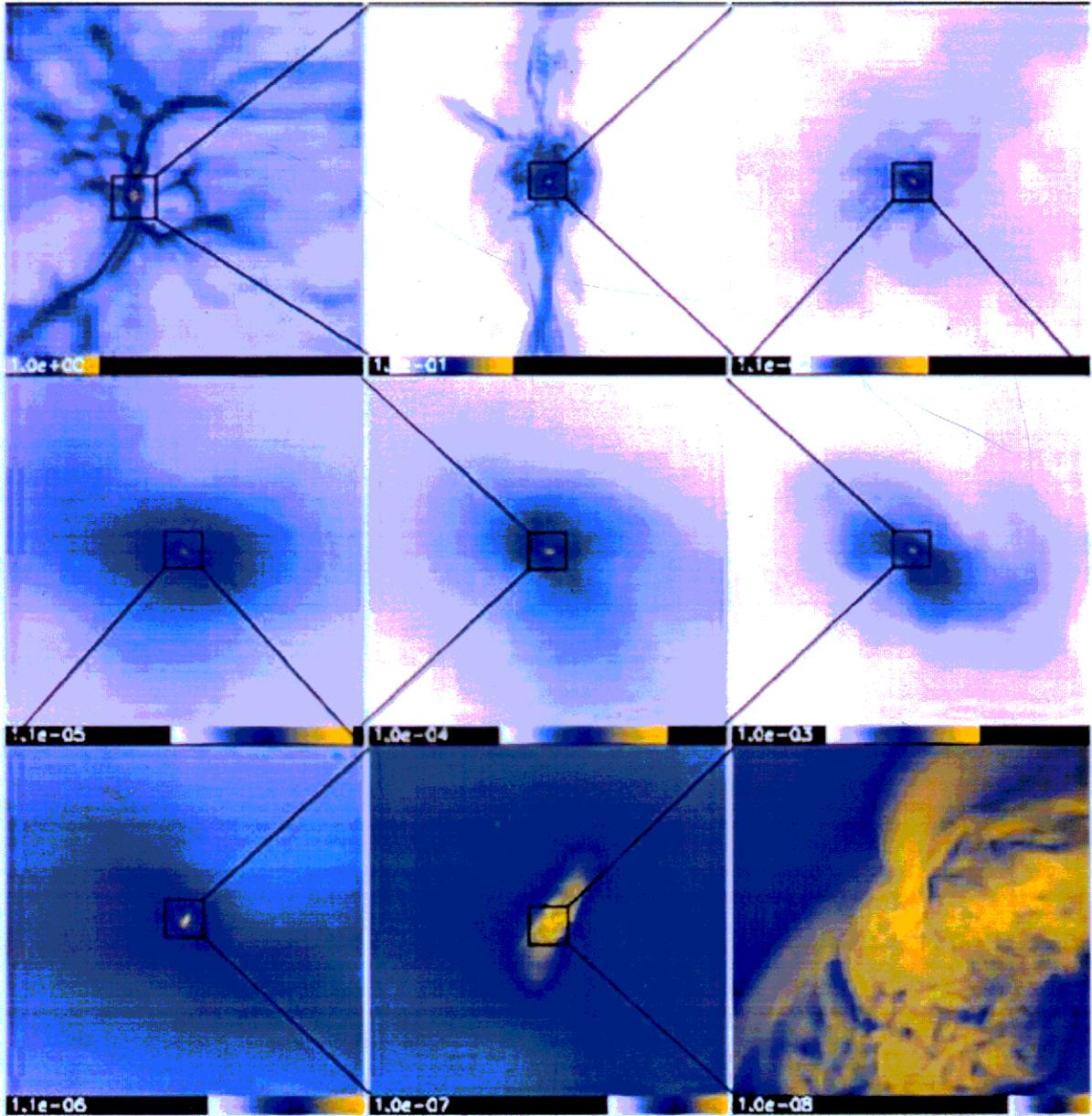
Fig. 3. Million-fold zoom showing the multiscale structure of a primordial protostellar cloud core in the center of a low mass halo at $z=19.1$. Plotted is the logarithm of the baryon overdensity on a slice passing through the densest structure on the grid. Zoom proceeds clockwise from upper left. Linear scales are proper. The smallest grid cell in the center of the cloud core in the last frame is 2^{-4} pc = 43 AU.

(NORMAN, ABEL, & BRYAN 2000)

FORMATION OF $\sim 200 M_{\odot}$ CORE $\rightarrow \sim 100 M_{\odot}$ STAR?

LARSON (1999):

- FROM THEORY: AT $z=0$: JEANS MASS $\sim 1000 M_{\odot} \rightarrow 100 \dots 500 M_{\odot}$ STARS
 \rightarrow IMF DOMINATED BY MASSIVE & VERY MASSIVE STARS
- DUE TO MUCH HIGHER TEMPERATURE
 - \leftrightarrow COOLING ONLY BY TRACES OF MOLECULAR HYDROGEN
 - (SINCE THERE ARE NO METALS)



ABEL, BRYAN, NORMAN (2001)

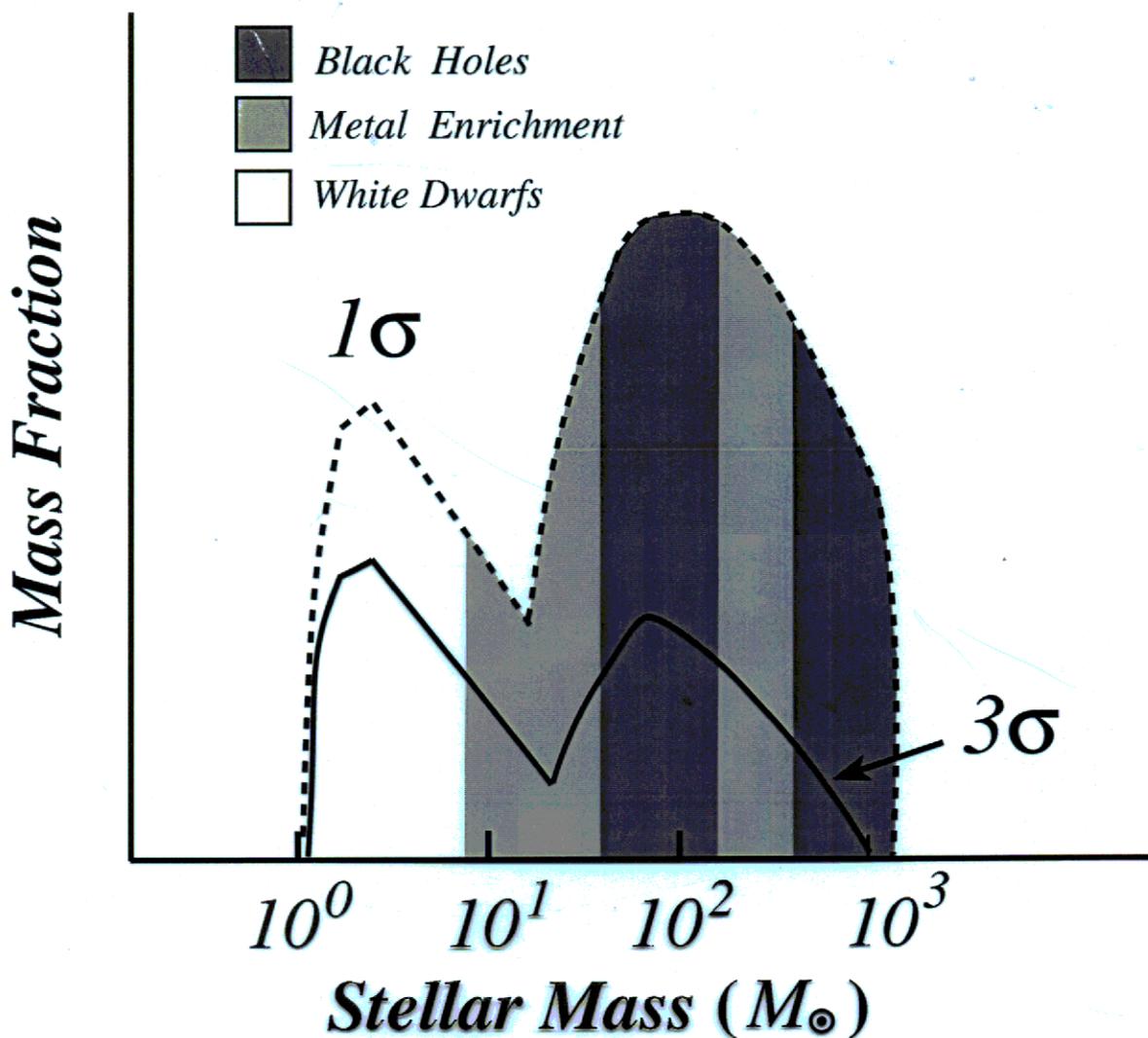
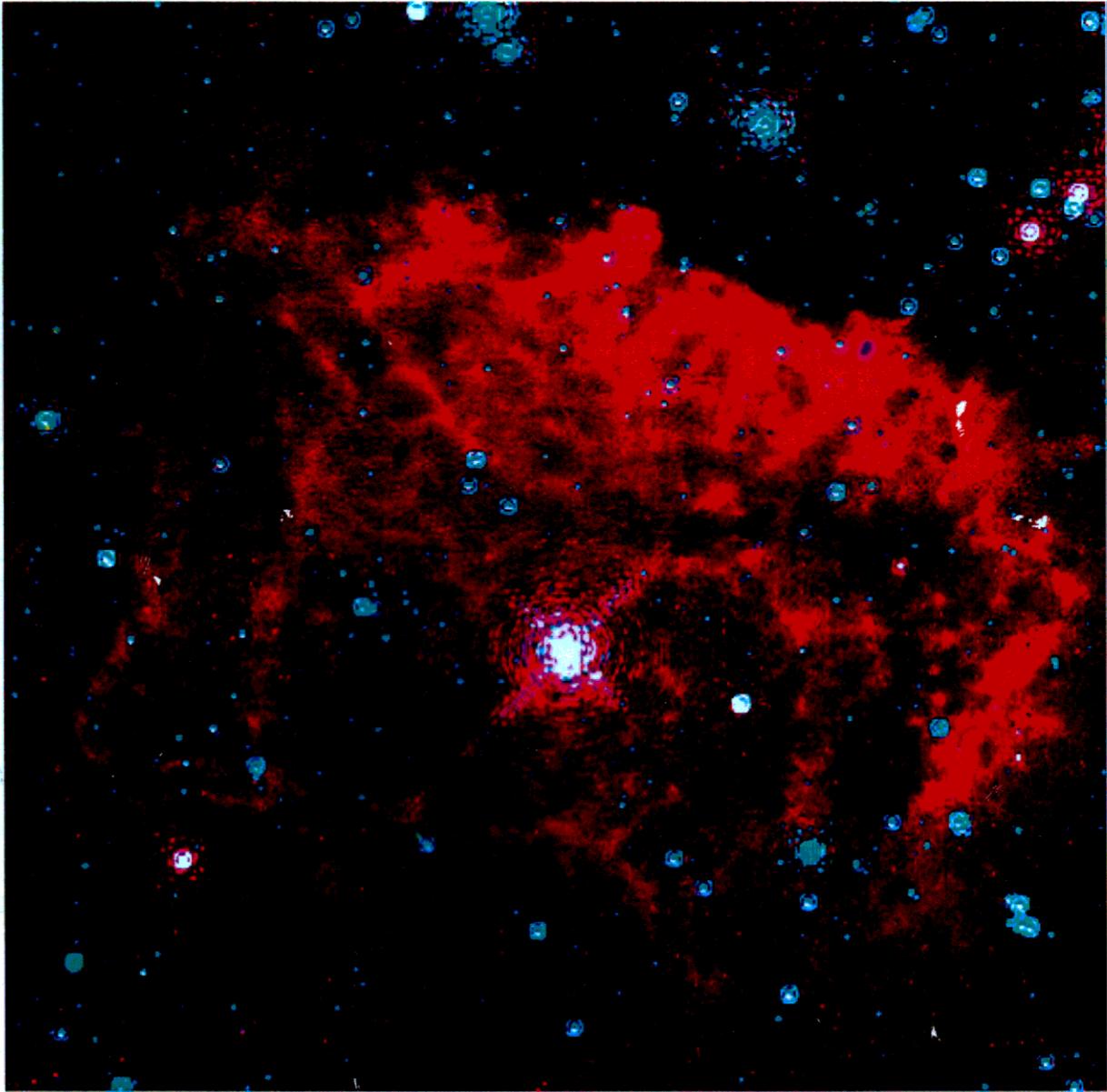


FIG. 8.—Schematic IMF of Population III stars. Solid and dashed lines are the IMFs for 3σ and 1σ density fluctuations, respectively. The IMF of Population III stars is likely to be bimodal and approximated by a superposition of two power-law-like components with two different peaks of $m_{p1} \approx 1-2 M_{\odot}$ and $m_{p2} \approx$ a few times $10-10^2 M_{\odot}$. The relative height of the first peak descends with time compared to the second peak. According to a recent theory of stellar evolution, stars with masses between $1-2$ and $8 M_{\odot}$ are likely to evolve to white dwarfs that may reside in galactic halos as baryonic dark matter. Stars with masses between 8 and $35 M_{\odot}$ probably evolve to supernovae and eject heavy elements in the intergalactic medium. Stars with masses between 35 and $10^2 M_{\odot}$ and greater than $250 M_{\odot}$ are likely to collapse into black holes and may be responsible for baryonic dark matter. Stars with masses between 10^2 and $250 M_{\odot}$ probably explode as supernovae and inject heavy elements into the intergalactic medium.

THE PROBLEM

CAN VERY MASSIVE STARS RETAIN THEIR MASS ?



PRC97-33 • ST ScI OPO • October 8, 1997 • D. Figer (UCLA) and NASA

KUDRITZKI (2000):

RADIATION-DRIVEN WINDS:

$$\dot{M} \sim \sqrt{Z}$$

(OVER WIDE RANGE OF Z)

LANERS ET AL. (2001):

$$\dot{M} \sim Z^{0.6}$$

THE PISTOL STAR

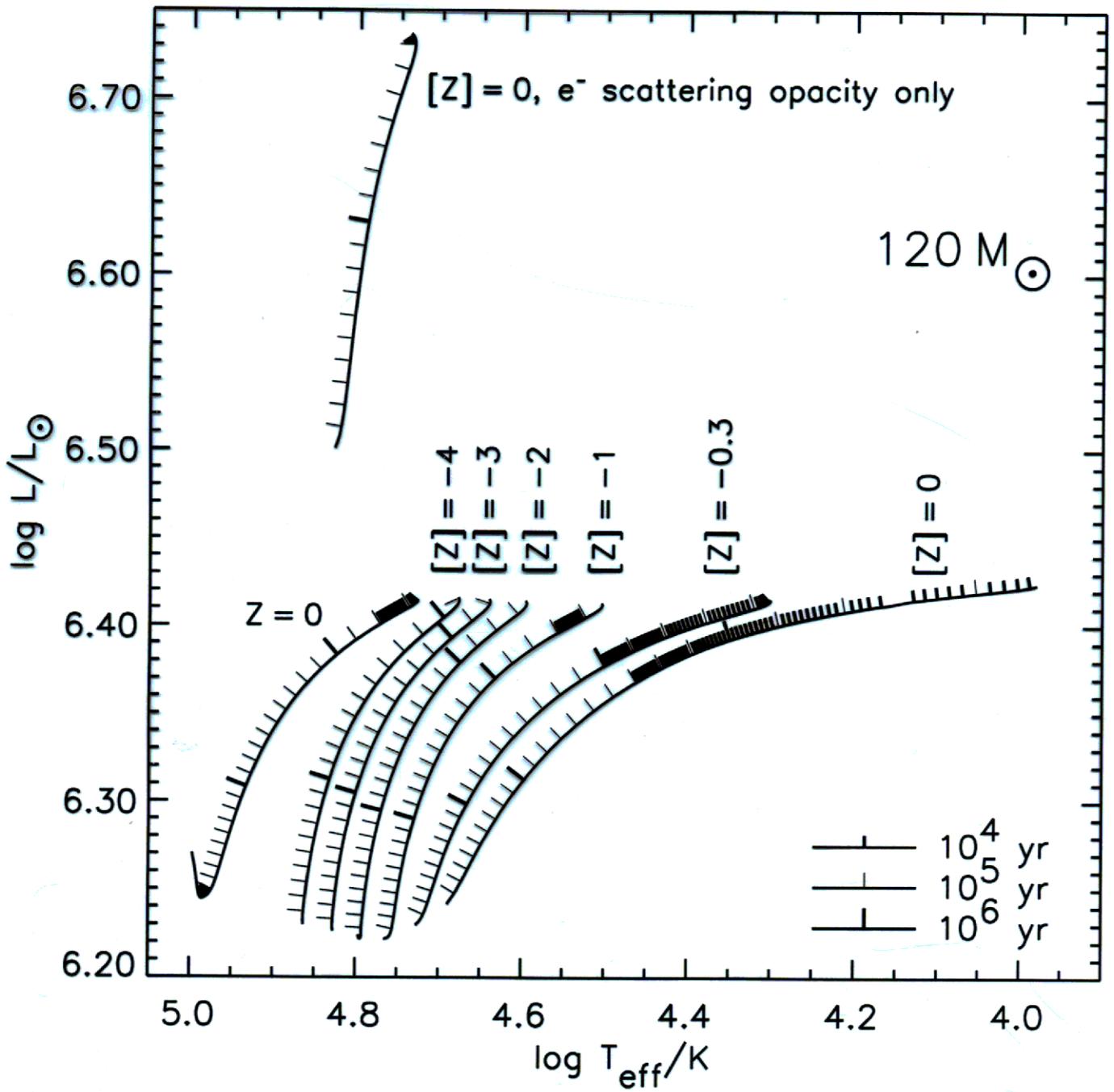
$\sim 150 \dots 200 M_{\odot}$

GALACTIC STAR
EXTREMELY HIGH
MASS LOSS RATE

UNKNOWN: CONTINUUM-DRIVEN WINDS

Mass Loss in Very Massive Primordial Stars

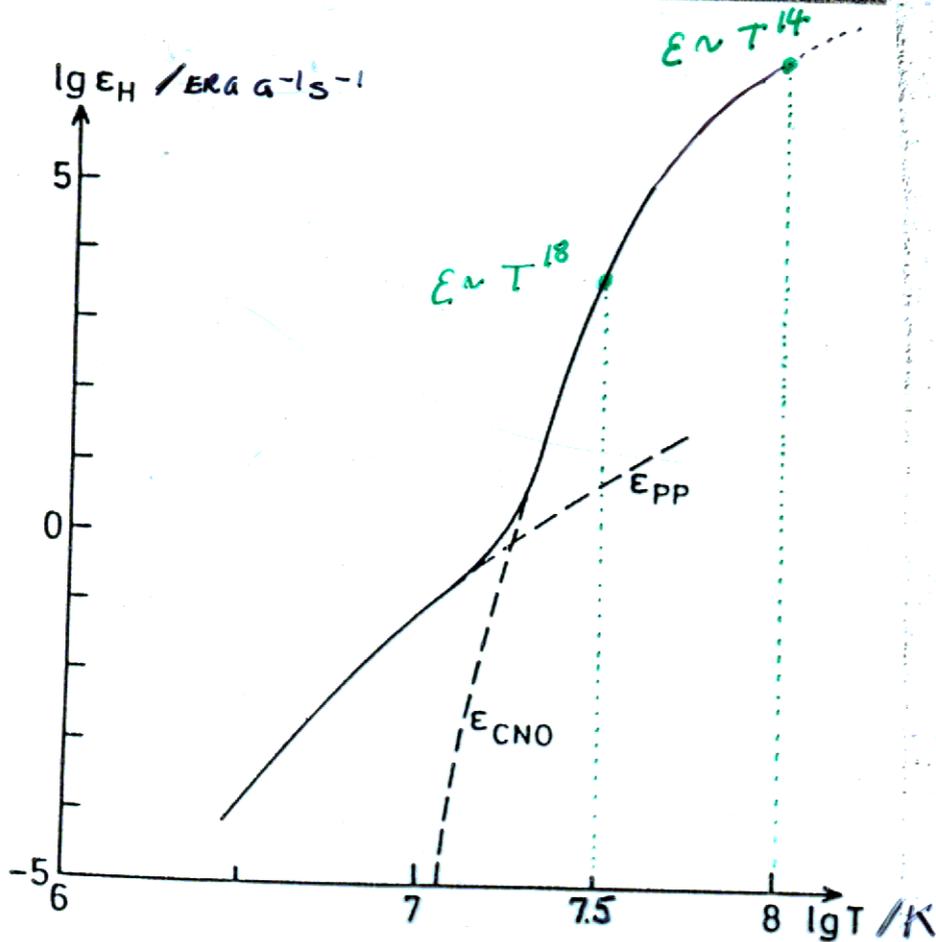
- negligible line-driven winds
(mass loss $\propto \sqrt{\text{metallicity}}$)
- no opacity-driven pulsations
(no metals)
- continuum-driven winds not yet understood
likely small contribution
- epsilon mechanism inefficient in metal-free stars
below $\lesssim 1000$ solar masses
From pulsational analysis we estimate:
 - 120 solar masses: $< 0.2\%$
 - 300 solar masses: $< 3.0\%$
 - 500 solar masses: $< 5.0\%$
 - 1000 solar masses: $< 12\%$during central hydrogen burning.
- **Red Super Giant** pulsations could lead to significant mass loss during helium burning for stars $\gtrsim 500$ solar masses



KUDRITZKI (2000): RADIATION DRIVEN WINDS: $M \sim \sqrt{Z}$

→ WHAT OTHER MASS LOSS MECHANISMS EXIST?

CALCULATIONS:
SHOW THAT THIS
RELATION SEEMS
VALID DOWN TO
QUITE LOW
VALUE OF Z !



THE ϵ -MECHANISM

[ϵ := SPECIFIC ENERGY GENERATION RATE]

FOR DRIVING PULSATIONS

$\frac{\partial \ln \epsilon}{\partial \ln T} \gg 1$, $\frac{d \ln L}{d \ln M} > 1 \rightarrow$ EFFECT BECOMES MORE IMPORTANT FOR HIGHER MASS

$\rightarrow \epsilon \Big|_{T_{max}}^{S_{max}} \gg \bar{\epsilon} \approx \frac{S_{max}}{S} \rightarrow$ DRIVING OF PULSATIONS

+ DARPENING IN ENVELOPE

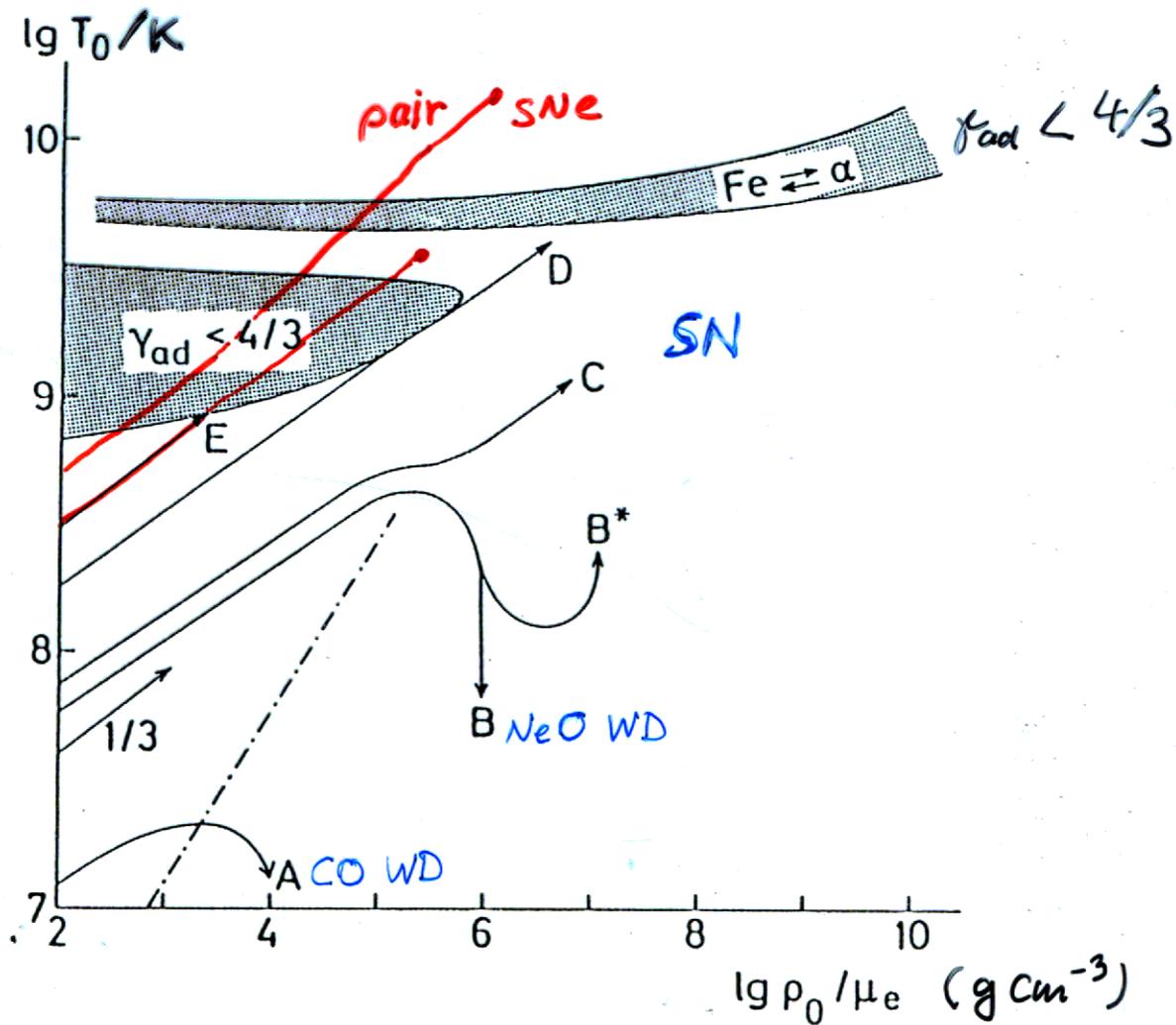
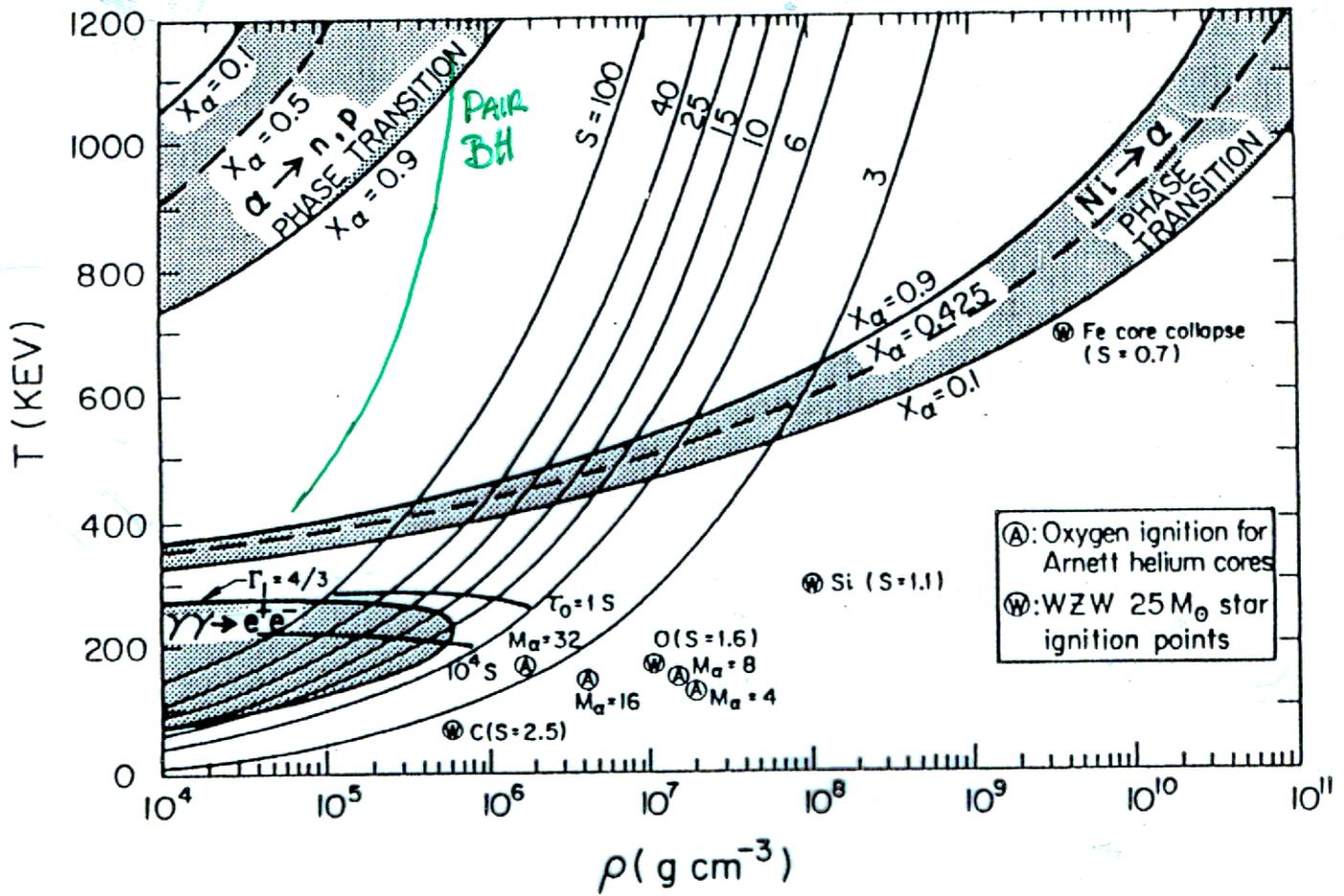
AT $Z=0$

$\frac{\partial \ln \epsilon}{\partial \ln T} \downarrow$
DARPENING IN ENVELOPE \downarrow
} $\rightarrow ?$

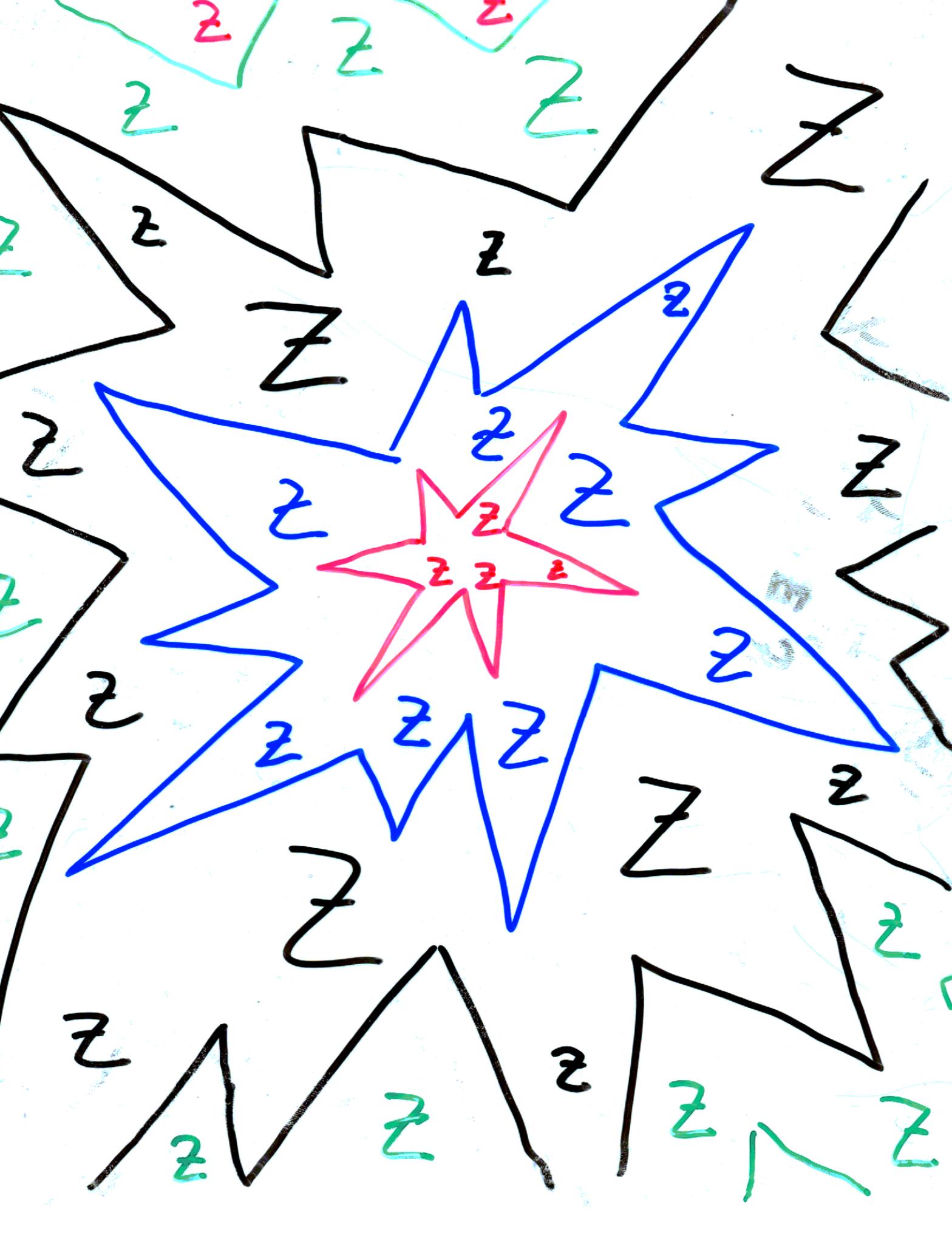
← HIGHER T
SINCE NO INITIAL CNO
 $\rightarrow T_c \approx 10^8 K$ 3d PROCESS
PRODUCES PRIMARY C

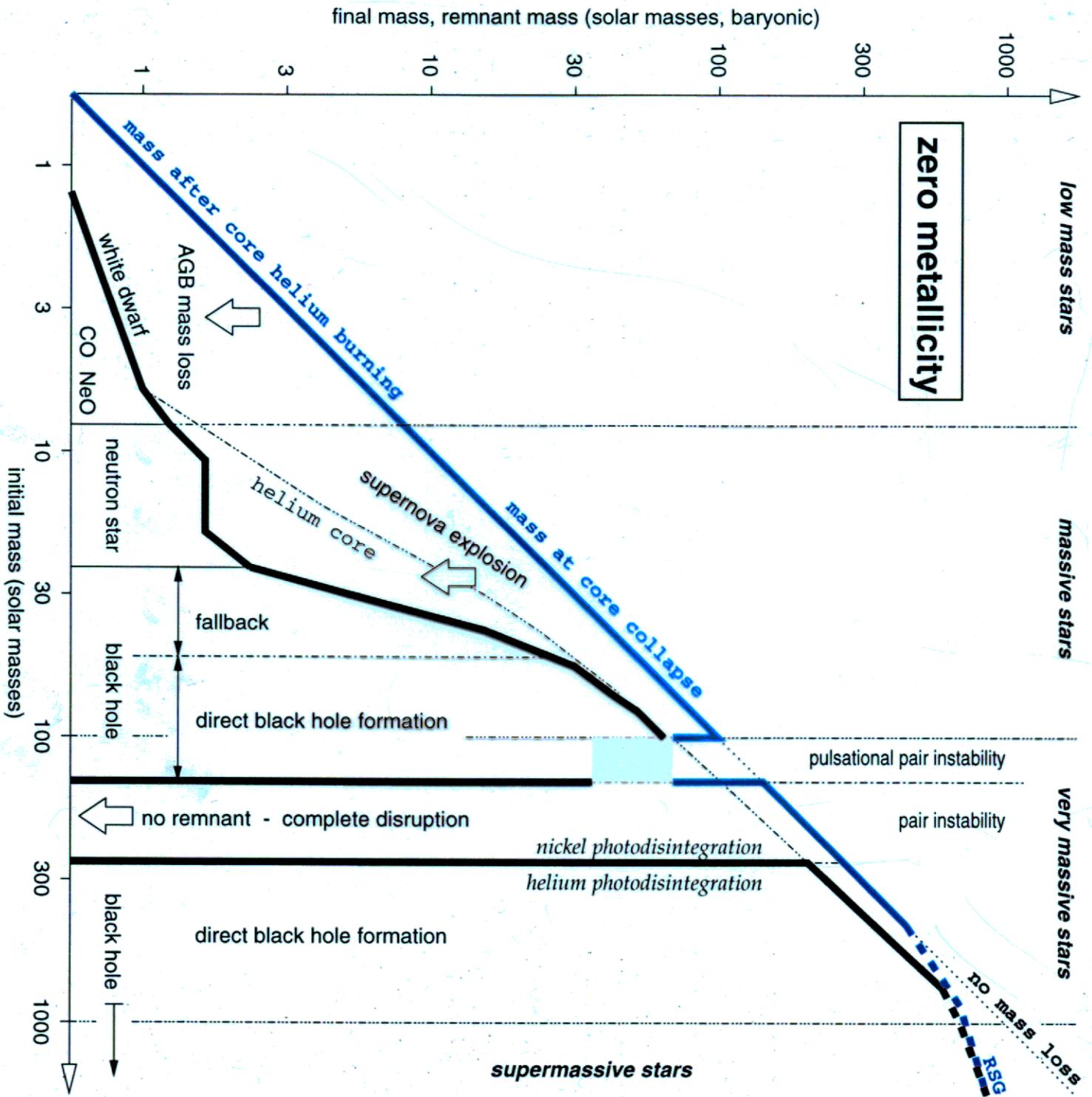
← MORE COMPACT

BOND, ARNETT, AND CARR (1984)

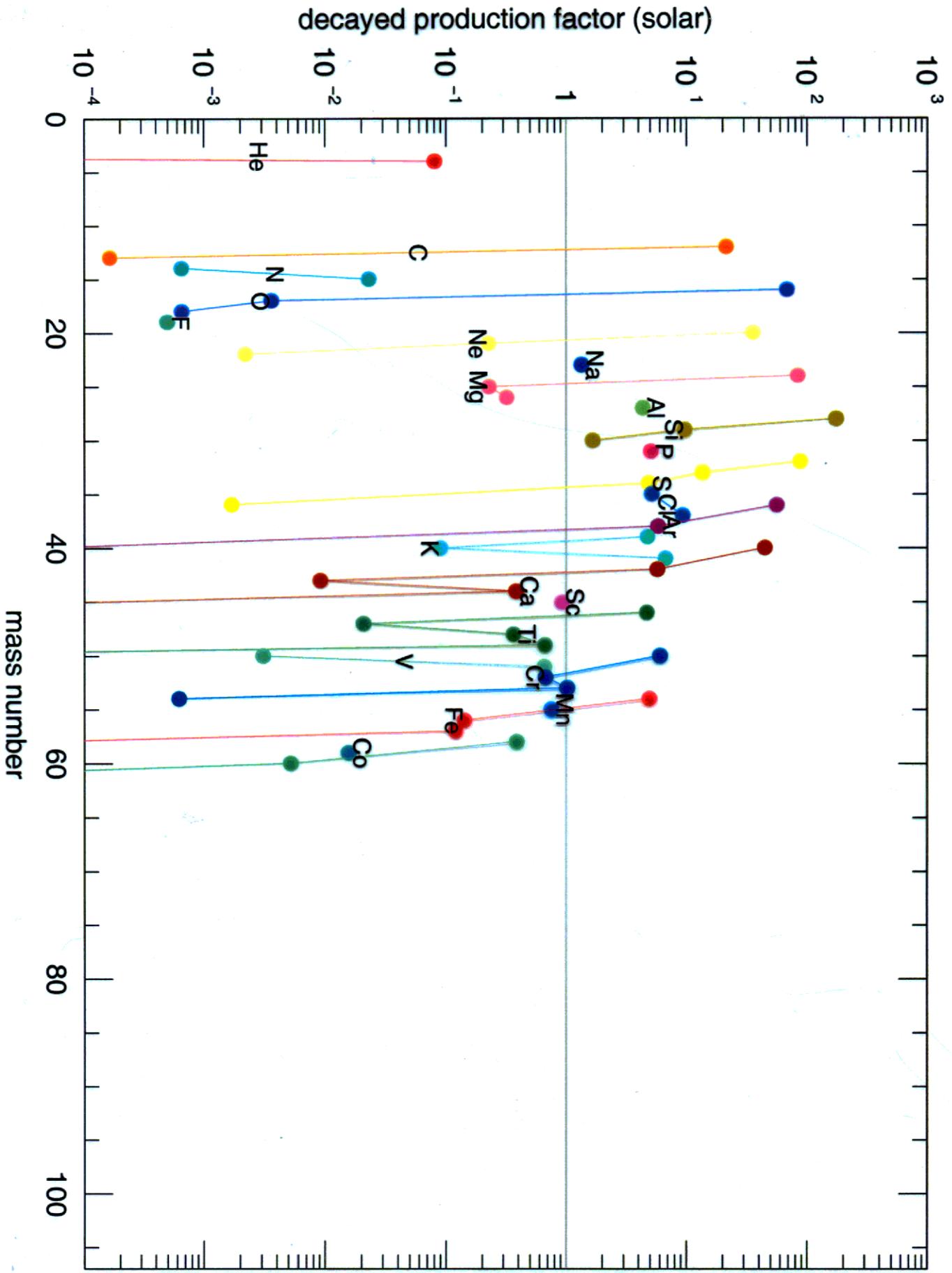


(KIPPENHAHN & WEIGERT, 1992)



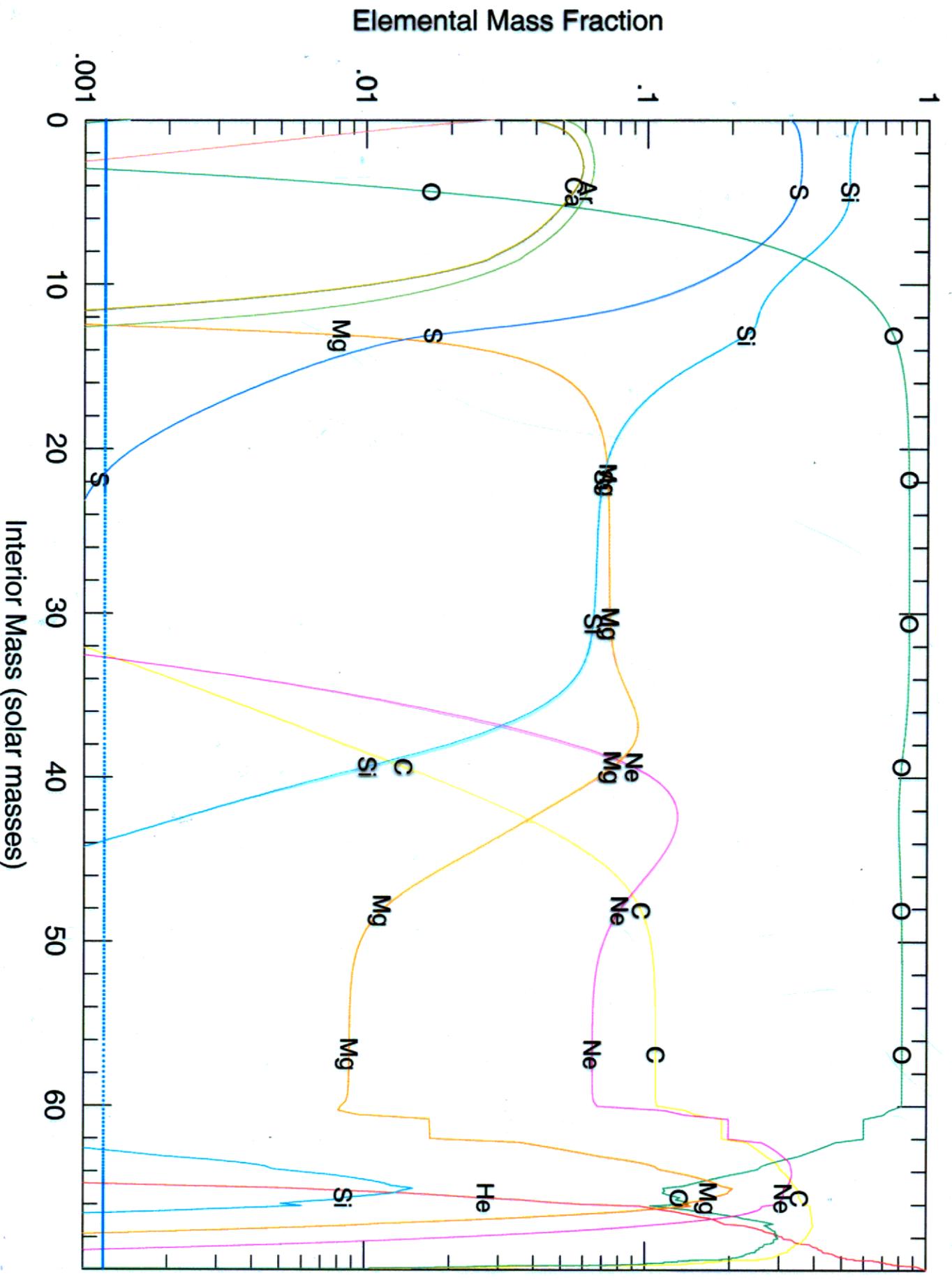


70 solar mass He core, primordial composition

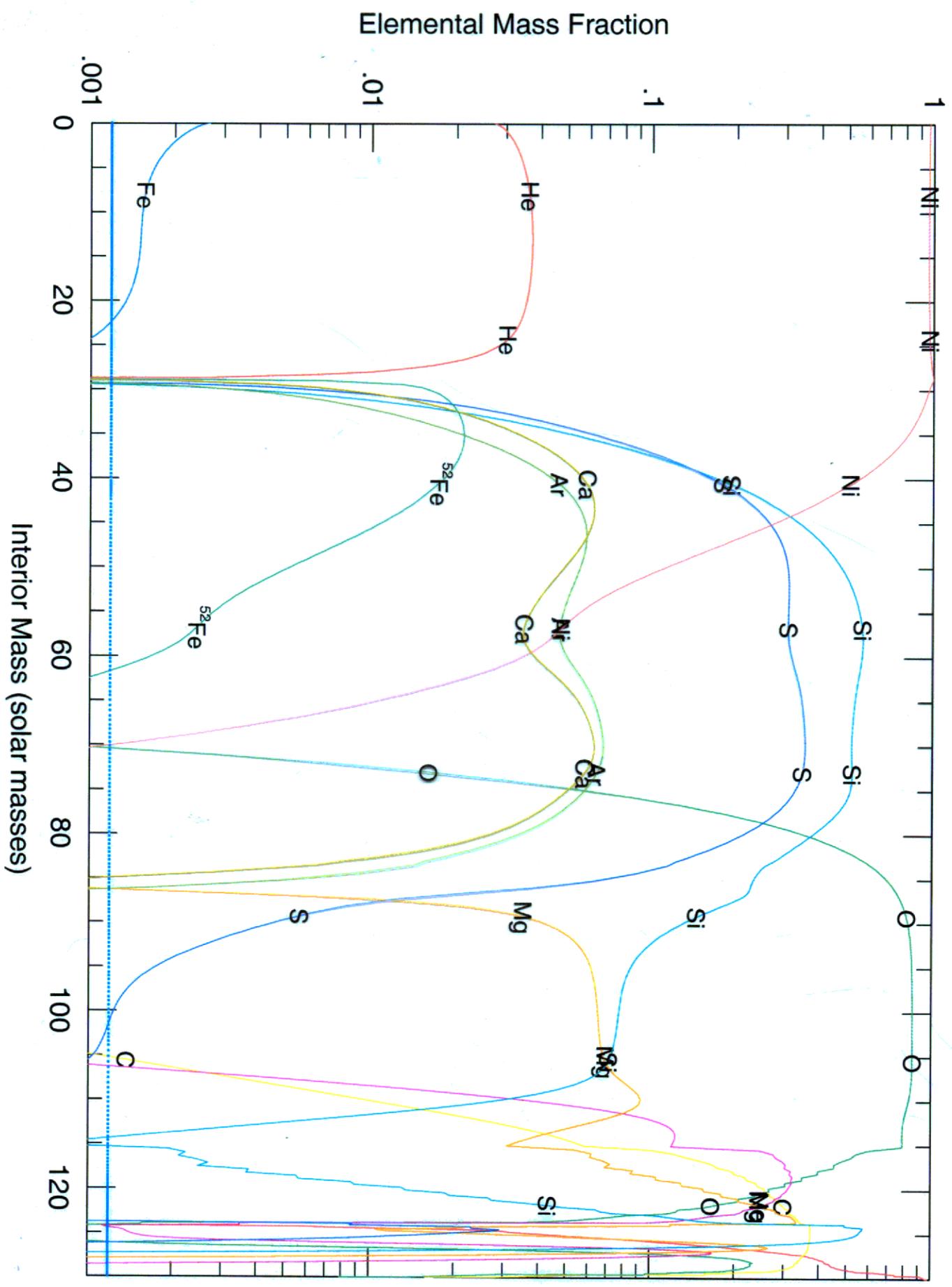


NO SOLID LINES IN CORE, PRIMORDIAL COMPOSITION

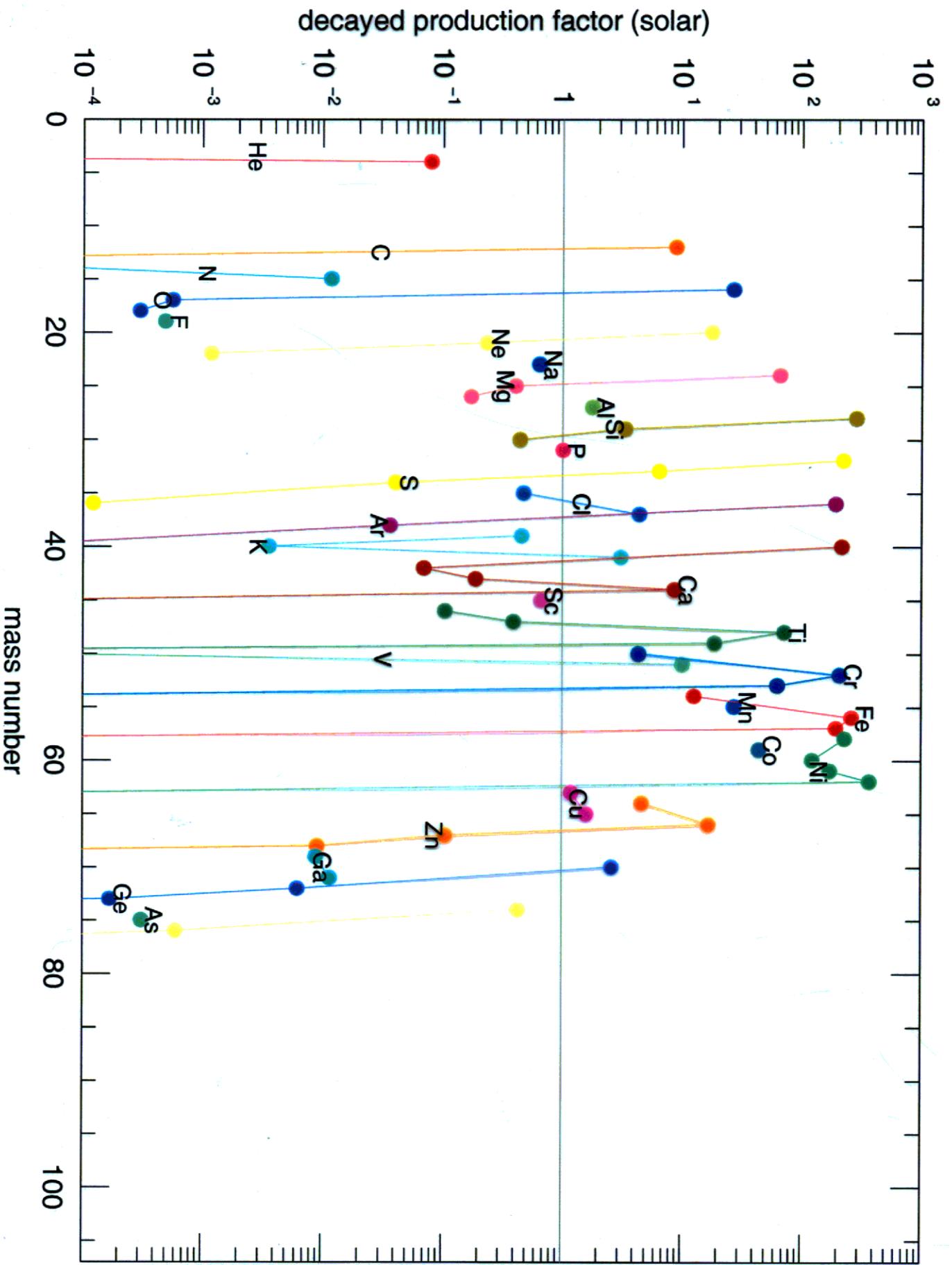
~ 160 M_☉ Star

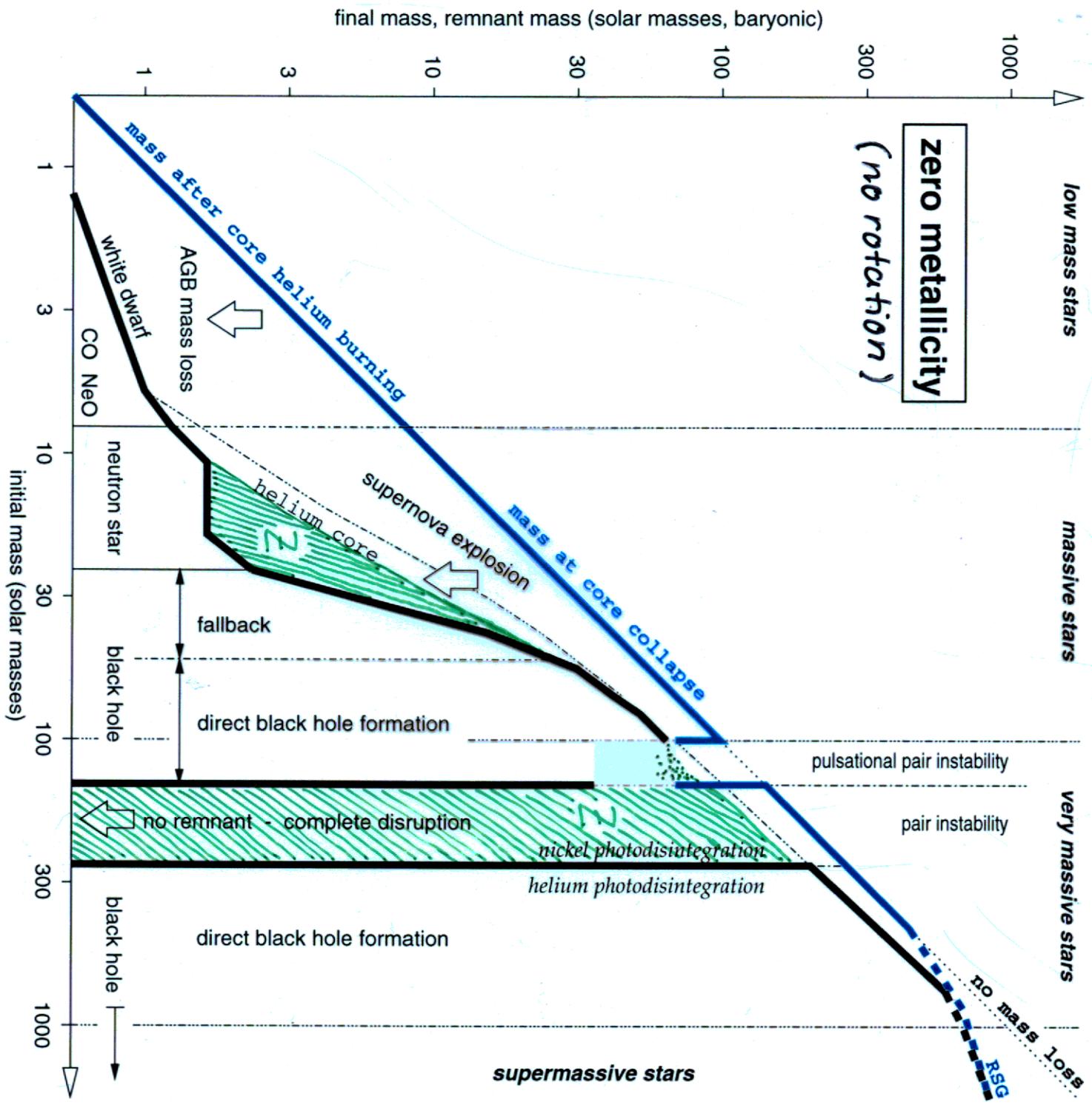


130 solar mass He core, primordial composition
 $\sim 250 M_{\odot}$ star



100 solar mass He core, primordial composition
~ 250 M_☉ star

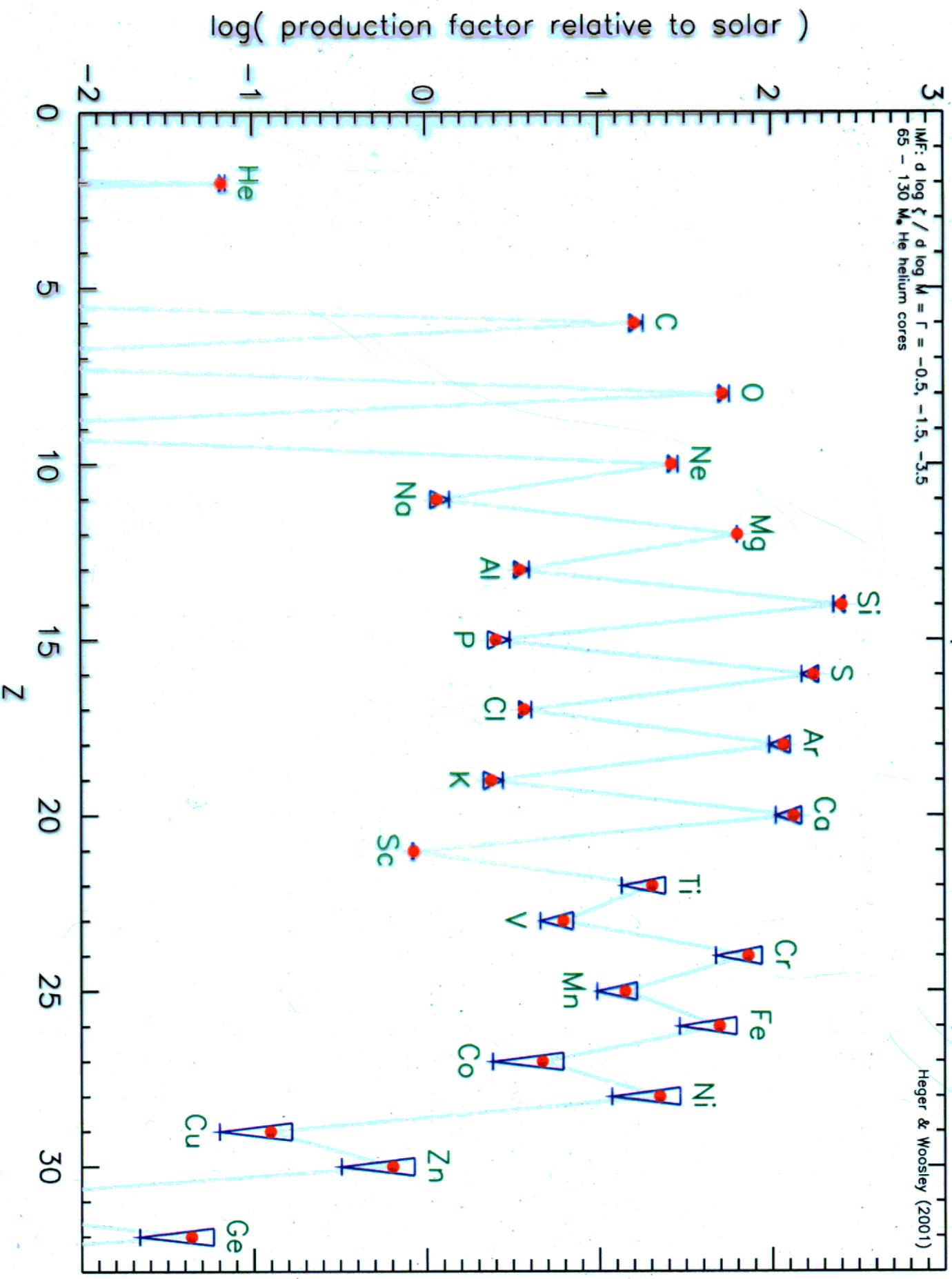




Production Factor of Pop III Pair Creation Supernovae

IMF: $d \log \xi / d \log M = \Gamma = -0.5, -1.5, -3.5$
65 - 130 M_{\odot} He helium cores

Heger & Woosley (2001)

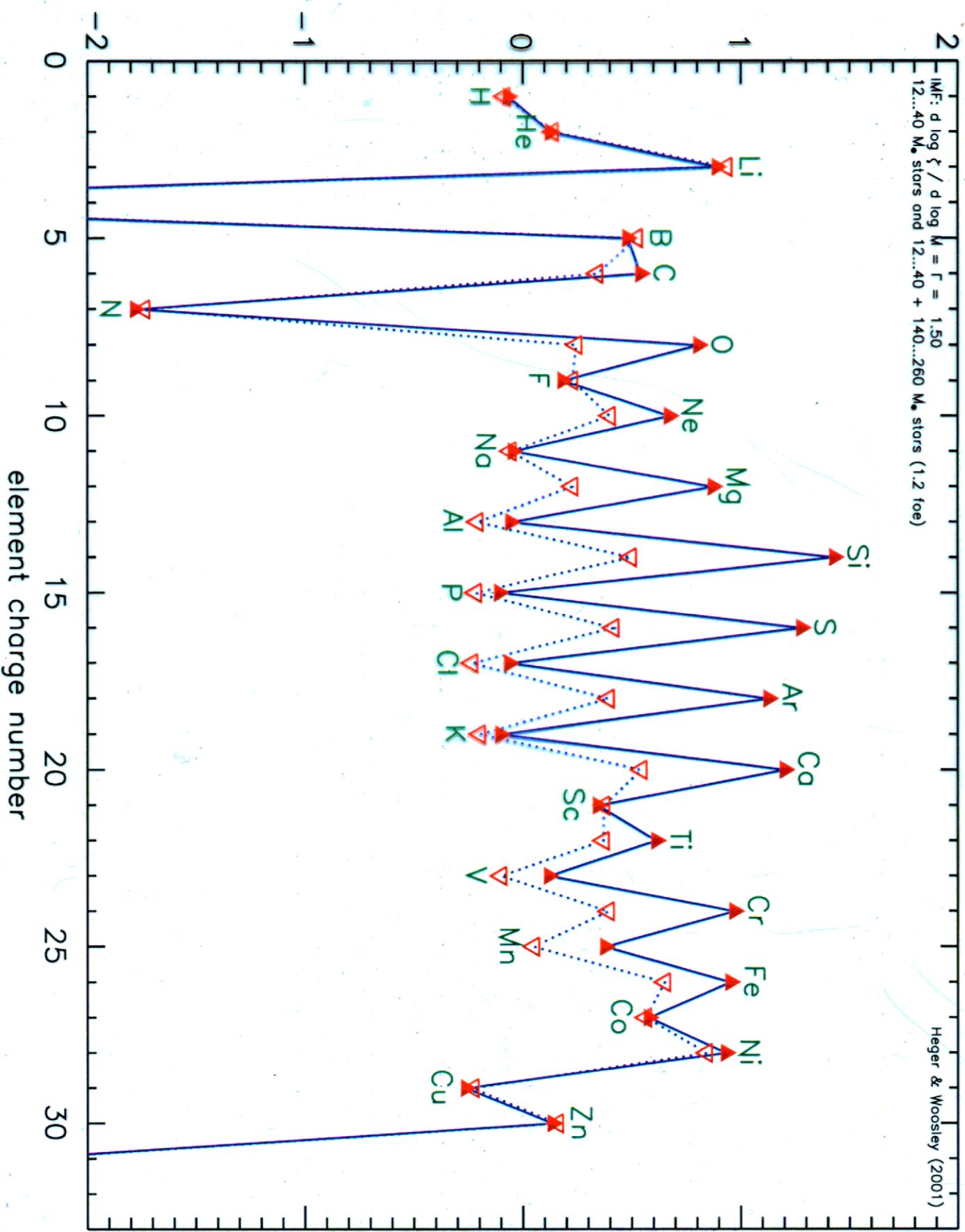


Production Factor of Primordial Massive Stars

IMF: $d \log \xi / d \log M = \Gamma = 1.50$
 12...40 M_{\odot} stars and 12...40 + 140...260 M_{\odot} stars (1.2 foe)

Heger & Woosley (2001)

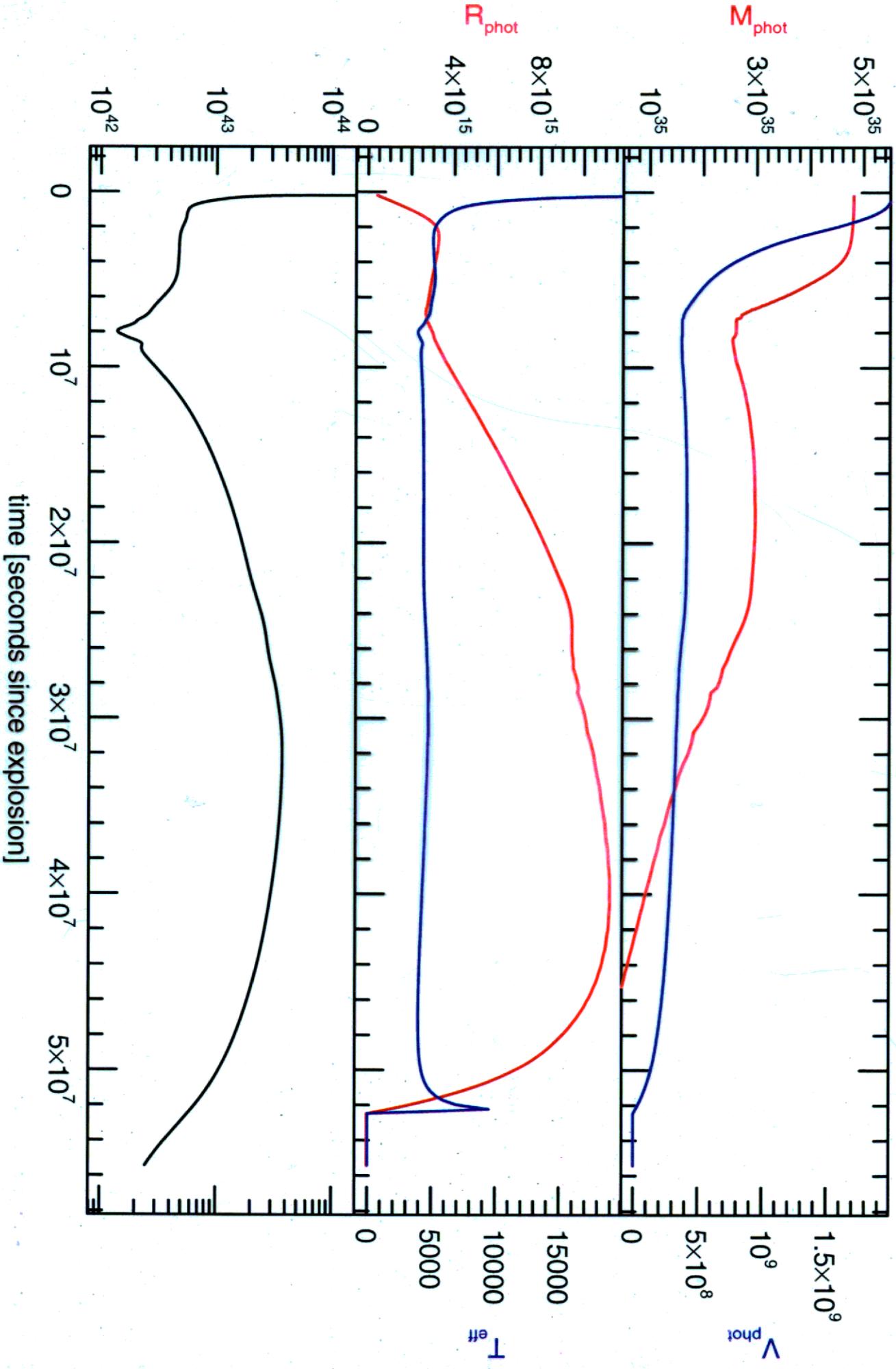
log(production factor relative to solar)



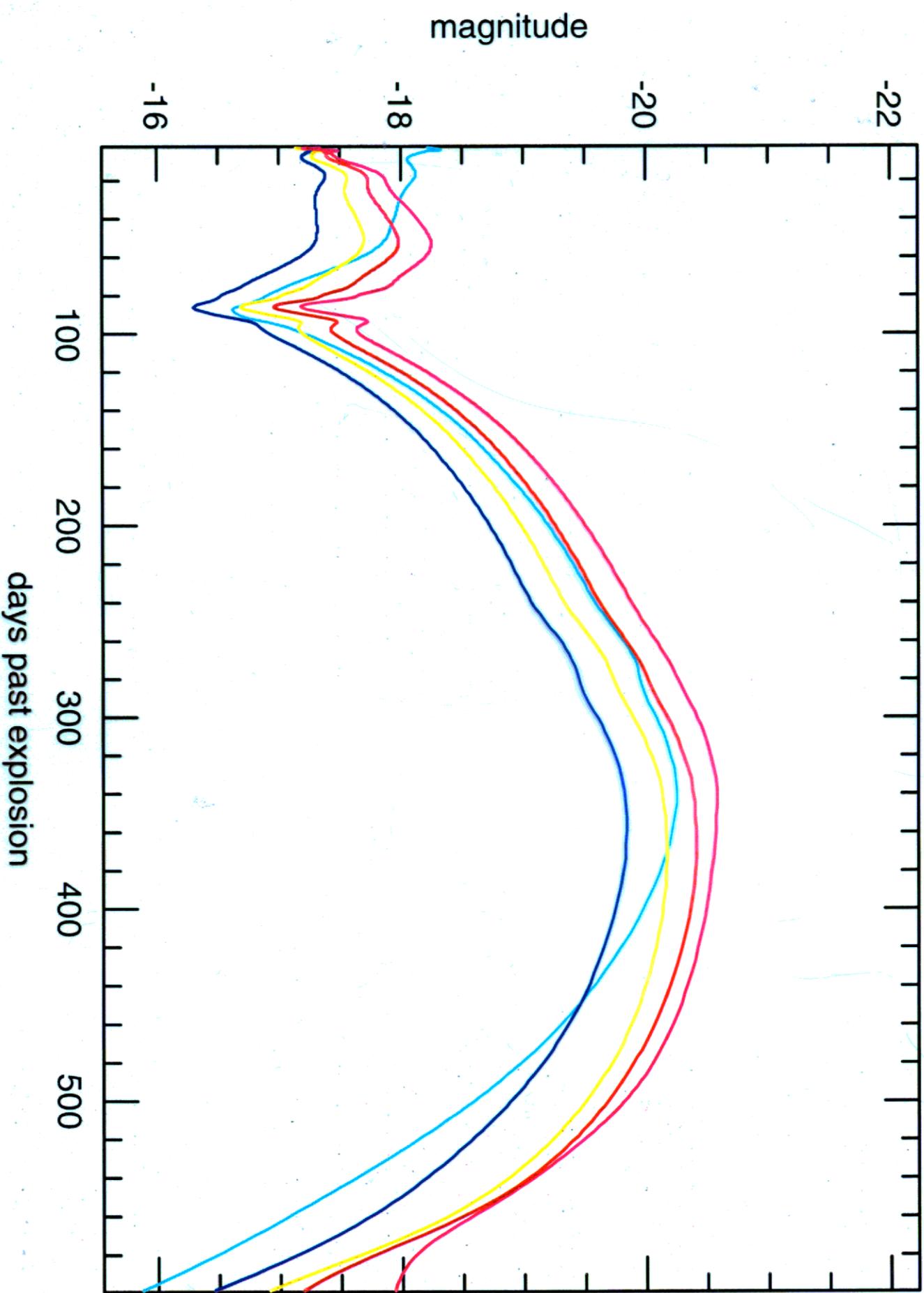
Bright Supernovae at the Edge of the Universe

- explosion energy up to 10^{53} erg
(50 – 100× that of “normal” supernovae)
- up to 50 solar masses of radioactive ^{56}Ni
(100× that of Type Ia supernovae)
- Assuming that 10^{-6} of all baryons go into stars of 250 solar masses, at a red shift of 20, the rate of events could be as high as **one every 6 sec.**
(For a current standard cosmology, $\Omega_\Lambda = 0.7$, $\Omega_{\text{matter}} = 0.3$, $H_0 = 65 \text{ km/s/Mpc}$, $\Omega_{\text{baryons}} = 0.02/h^2 = 0.047$)
- They would last about 10 yr in the observer frame
(due to large mass)
- This computes to ~ 1000 of these objects are visible per square degree at any given time
(assuming no extinction)
- They are intrinsically a few times brighter than Type Ia supernovae
(in bolometric luminosity)
Note that the bolometric luminosity of an object at a redshift of 20 is not much less than that of an object with at redshift of a few, where Type Ia supernovae have already been observed!
- Only observable in the near infrared
(due to absorption by neutral hydrogen short of 1215 Å, redshifted by a factor $1 + \text{redshift}$)

PHIL PINTO (2002)

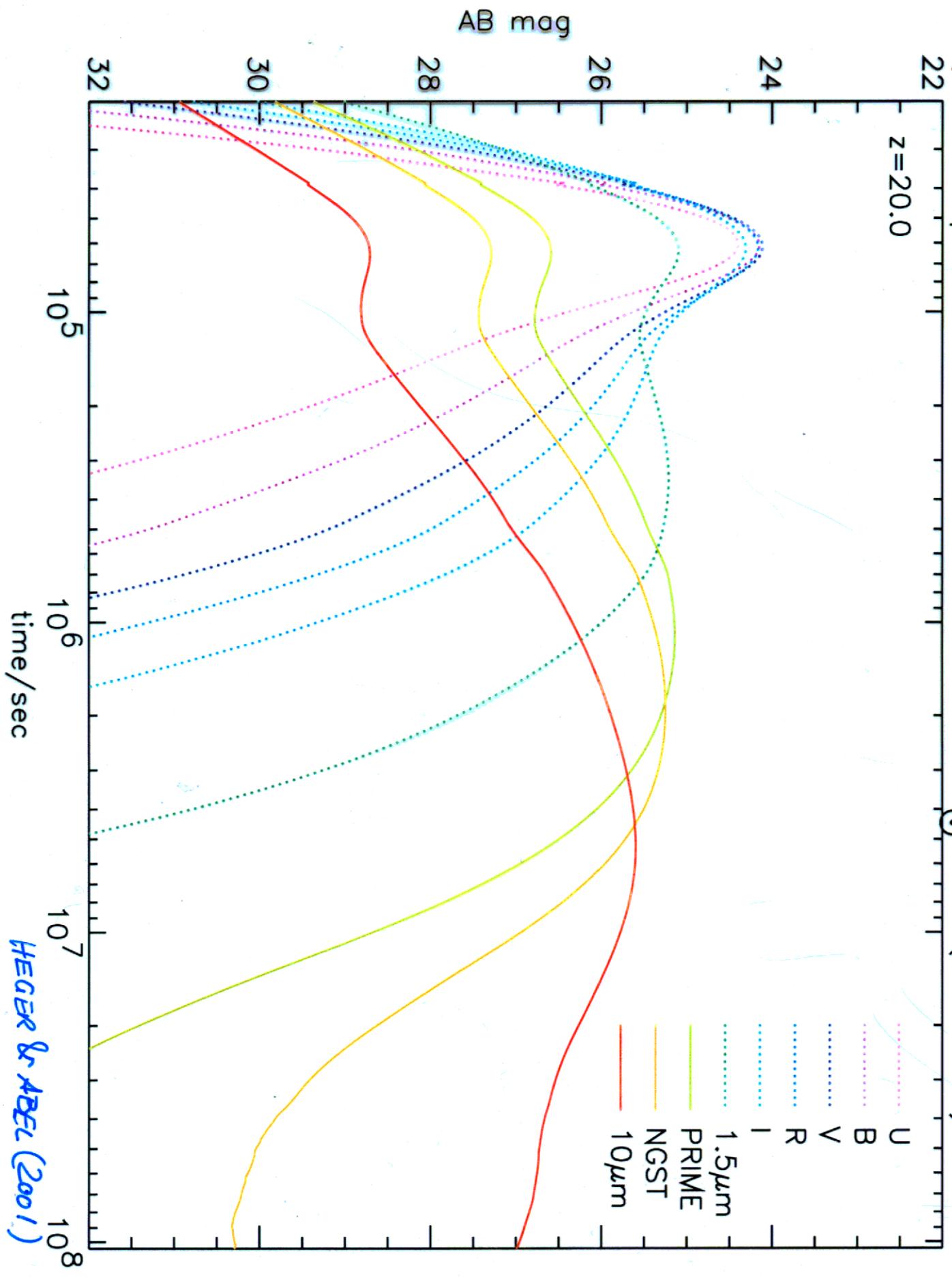


PHIL PINNOC (2002)



pair-SN of a metal-free 250 M_{\odot} star (64.5 foe)

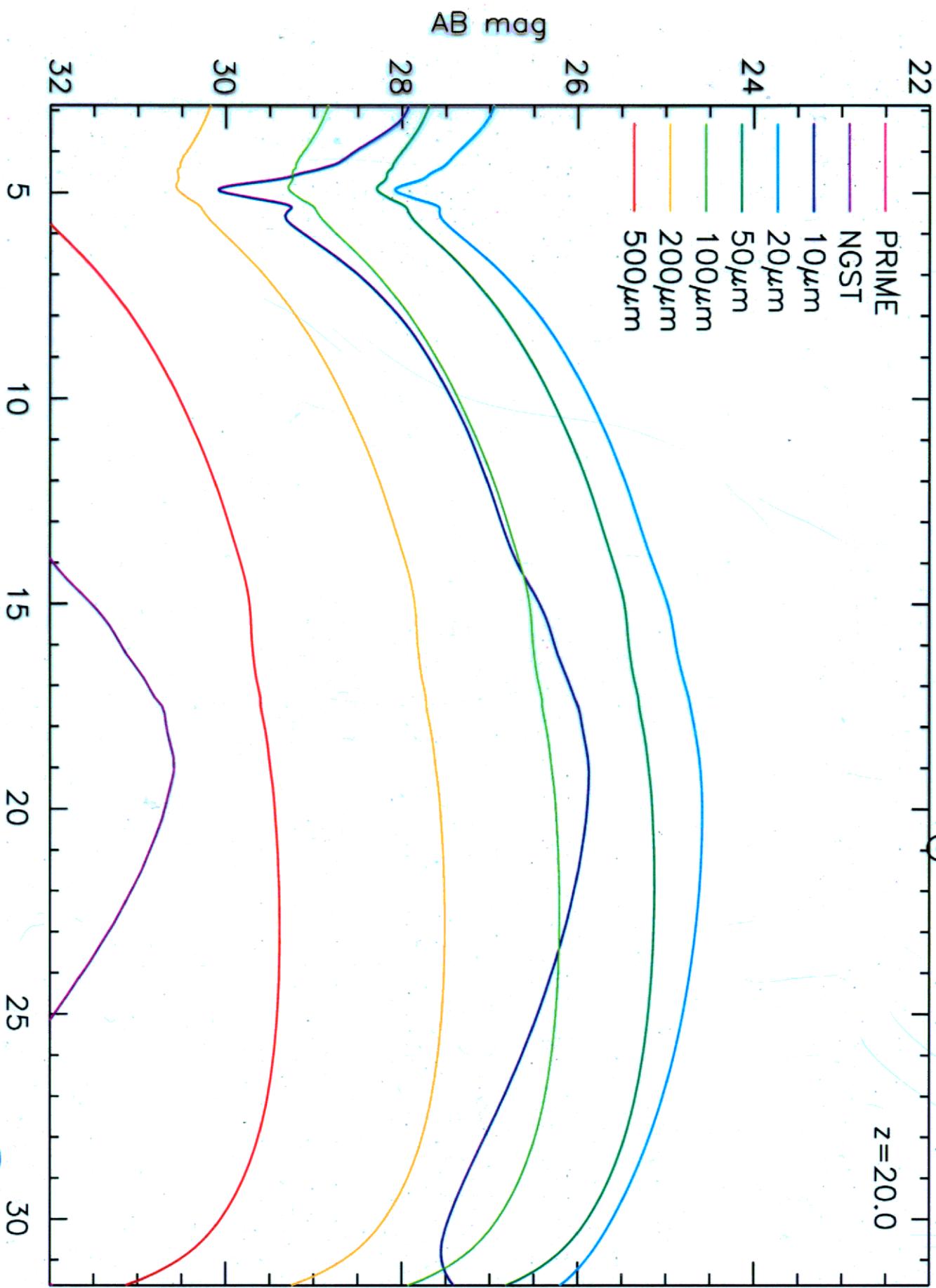
$z=20.0$



HEGER & ABEL (2001)

pair-SN of a metal-free 250 M_{\odot} star (64.5 foe)

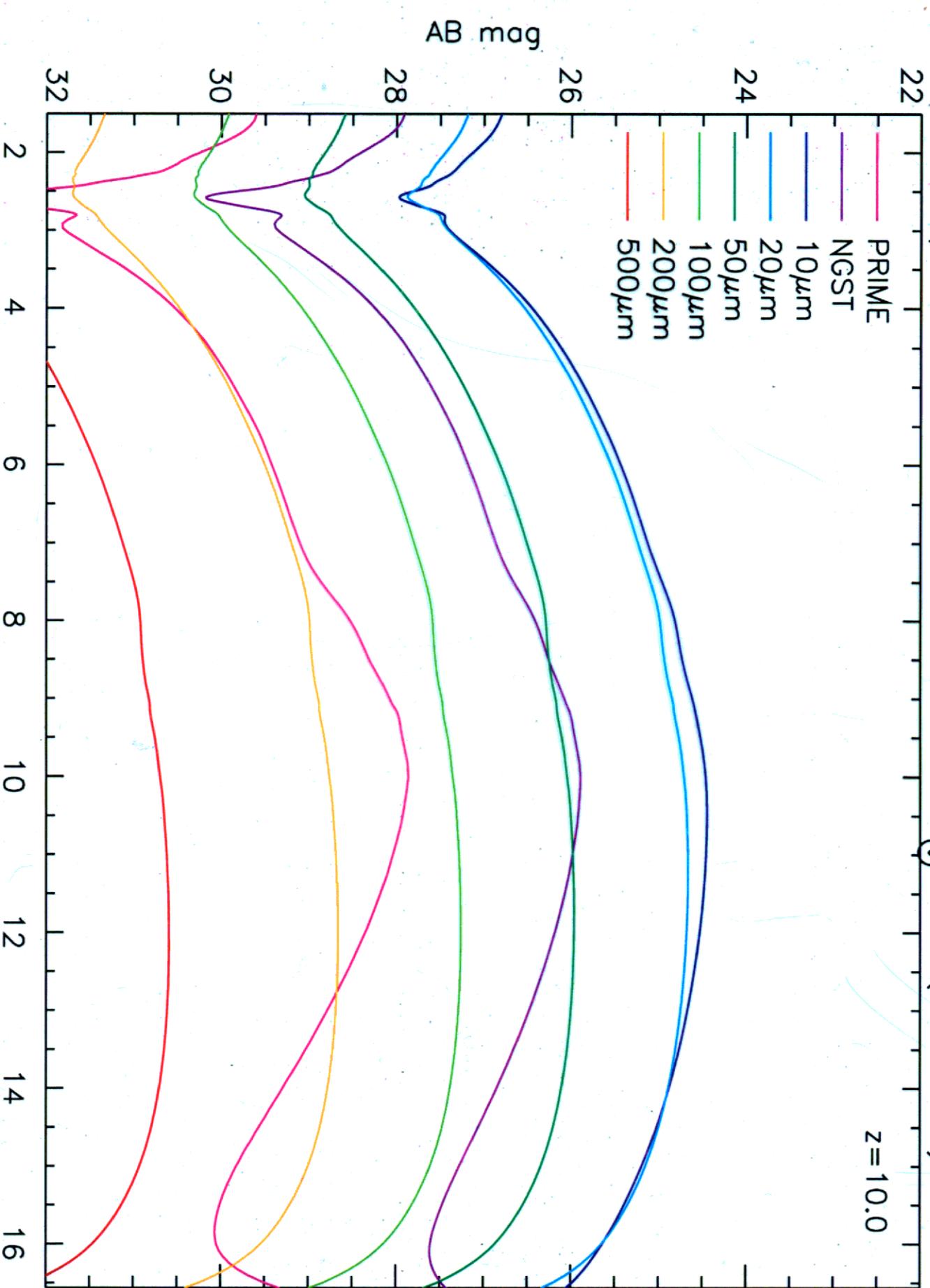
$z=20.0$



EPON LIGHT CURVE BY PHIL FINE (2002)

pair-SN of a metal-free 250 M_⊙ star (64.5 foe)

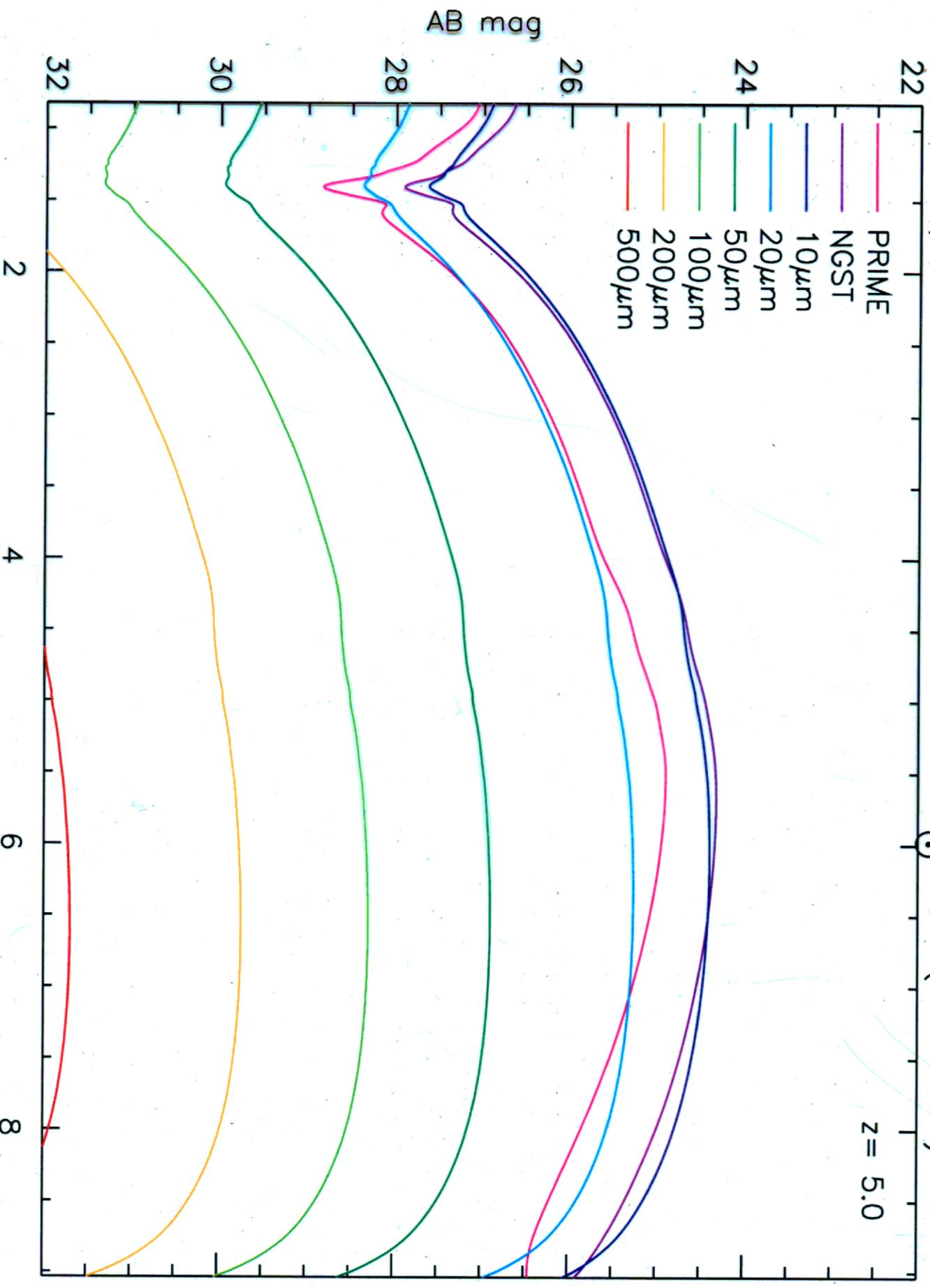
z=10.0



FRONT LIGHTCURVE BY PHIL PINTO (2002)

pair-SN of a metal-free 250 M_⊙ star (64.5 foe)

z = 5.0



END LIGHTCURVE BY PHIL PINTO

Summary

Due to their unique composition, the birth, life and death of the first stars is very different from later generations:

- They are likely born much more massive.
(mass scale: 100 solar masses)
- They are more compact and hotter.
(typically 100,000 K for a 100 solar mass star)
- Even stars of several 100 solar masses might not lose significant amounts of mass before death.
(no winds, no epsilon mechanism)
- They can encounter the pair-creation instability:
 - energetic explosions of up to 100× that of “normal” supernovae (10^{53} erg kinetic energy)
 - up to 50 solar masses ^{56}Ni ejected
 - might be seen directly by NGST – out to red shift 20!
- Peculiar nucleosynthesis pattern from pair-creation supernovae:
 - no heavy elements beyond zinc produced
 - no *r*-process, no *s*-process
- Odd-even elemental pattern in all massive stars, but particularly strong in pair-creation supernovae.

Outlook

- Influence of rotation on mass limits and nucleosynthesis.
- Multi-dimensional simulations of pair-instability supernovae.
- Detailed study of pulsational pair instability.
- Search for stars with VMS nucleosynthesis pattern.